



Double beta-decay workshop  
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## *Nucleon pairs in double-beta decay*

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Supported by MEXT and JICFuS as a priority issue (Elucidation of the fundamental laws and evolution of the universe) to be tackled by using Post 'K' Computer

## Outline

1. Neutrinoless double beta decay  $^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$
2. Decay from neutron closed-shell nuclei
3. Decay from spherical open-shell nuclei
4. Open question on rotational nuclei

# A large-scale shell model calculation on $^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$ neutronless double beta decay

Y. Iwata et al., PRL 116, 112502 (2016)

Among many earlier studies, we carried out somewhat larger shell-model calculations in the  $sd+pf$  shell.

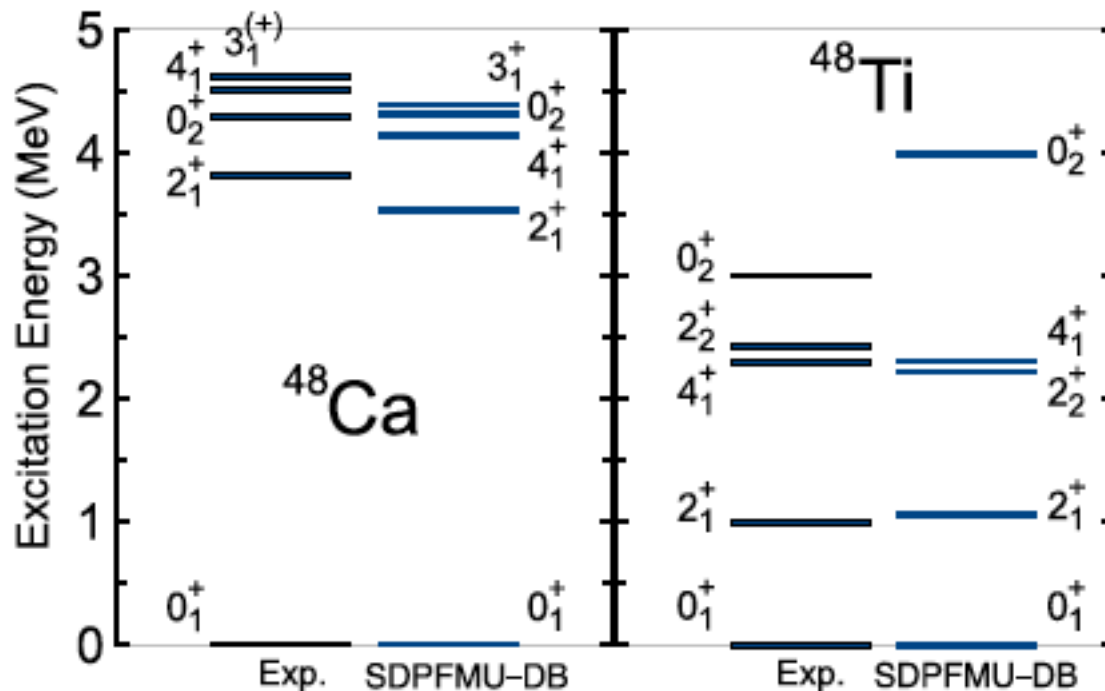


FIG. 1. Excitation spectra of  $^{48}\text{Ca}$  and  $^{48}\text{Ti}$ . The lowest five positive-parity states [41] are compared to  $sdpf$  calculations with the SDPFMU-DB interaction.

$$[T_{1/2}^{0\nu}(0_i^+ \rightarrow 0_f^+)]^{-1} = G^{0\nu} |M^{0\nu}|^2 \left( \frac{\langle m_{\beta\beta} \rangle}{m_e} \right)^2.$$

$$M^{0\nu} = \langle 0_f^+ | \hat{O}^{0\nu} | 0_i^+ \rangle = M_{\text{GT}}^{0\nu} - \frac{g_V^2}{g_A^2} M_{\text{F}}^{0\nu} + M_{\text{T}}^{0\nu}$$

TABLE I. NME value for the  $^{48}\text{Ca}$   $0\nu\beta\beta$  decay. The  $pf$ -shell calculation with GXPF1B is compared to the  $sdpf$   $2\hbar\omega$  results obtained with the SDPFMU-DB and SDPFMU interactions. Total values ( $M^{0\nu}$ ) are shown together with Gamow-Teller ( $M_{\text{GT}}^{0\nu}$ ), Fermi ( $M_{\text{F}}^{0\nu}$ ), and tensor ( $M_{\text{T}}^{0\nu}$ ) parts. Argonne- and CD-Bonn-type short-range correlations (SRC) are considered.

SRC	GXPF1B				SDPFMU-DB				SDPFMU			
	$M_{\text{GT}}^{0\nu}$	$M_{\text{F}}^{0\nu}$	$M_{\text{T}}^{0\nu}$	$M^{0\nu}$	$M_{\text{GT}}^{0\nu}$	$M_{\text{F}}^{0\nu}$	$M_{\text{T}}^{0\nu}$	$M^{0\nu}$	$M_{\text{GT}}^{0\nu}$	$M_{\text{F}}^{0\nu}$	$M_{\text{T}}^{0\nu}$	$M^{0\nu}$
None	0.776	-0.216	-0.077	0.833	0.997	-0.304	-0.067	1.118	0.894	-0.291	-0.068	1.007
CD-Bonn	0.809	-0.233	-0.074	0.880	1.045	-0.327	-0.065	1.183	0.939	-0.313	-0.065	1.068
Argonne	0.743	-0.213	-0.075	0.801	0.953	-0.300	-0.065	1.073	0.852	-0.288	-0.068	0.963

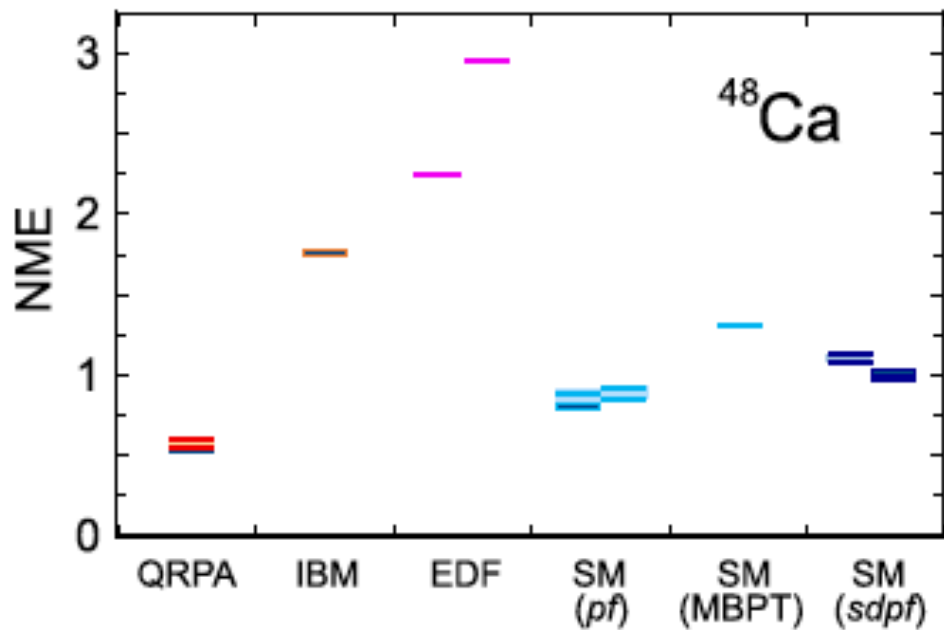
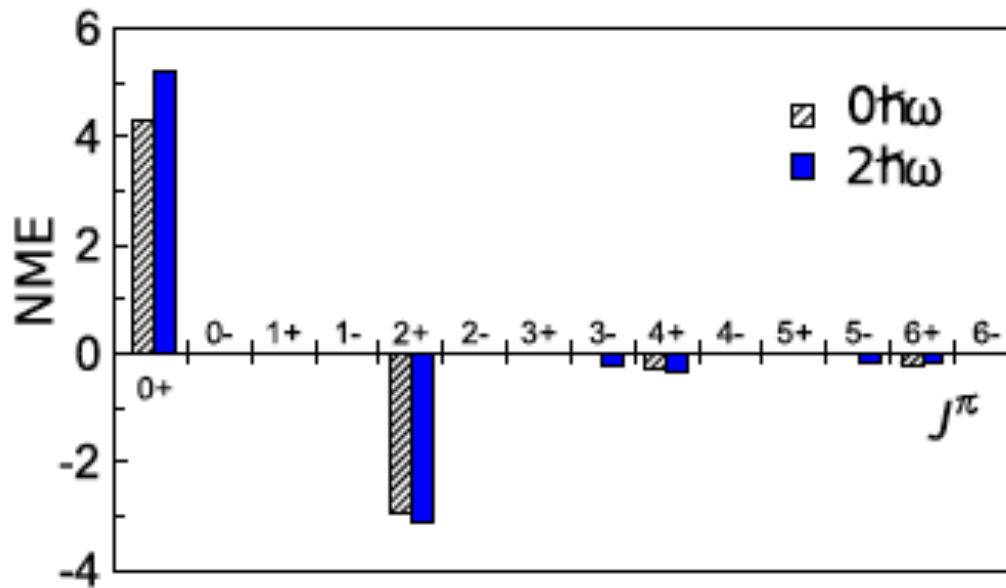


FIG. 2. Comparison of NME values for the  $^{48}\text{Ca}$   $0\nu\beta\beta$  decay. The present shell-model results in the  $sdpf$  space (SM  $sdpf$ : left SDPFMU-DB, right SDPFMU) are compared to  $pf$ -shell results (SM  $pf$ : left [17], right [15]),  $pf$ -shell result plus a perturbative calculation of the effect of orbitals outside the  $pf$  shell (SM MBPT) [50], QRPA [22], IBM [25], and EDF (left: nonrelativistic [26], right: relativistic [27]) calculations. The range between double horizontal bars covers results including a different type of short-range correlations (Argonne, CD-Bonn, UCOM [51]) and without them.

- SM (pf) [15] J. Menéndez, A. Poves, E. Caurier, and F. Nowacki, *Nucl. Phys. A* **818**, 139 (2009).
- [17] R. A. Sen'kov and M. Horoi, *Phys. Rev. C* **88**, 064312 (2013).
- MBPT [50] A. A. Kwiatkowski *et al.*, *Phys. Rev. C* **89**, 045502 (2014).
- QRPA [22] F. Šimkovic, V. Rodin, A. Faessler, and P. Vogel, *Phys. Rev. C* **87**, 045501 (2013).
- IBM [25] J. Barea, J. Kotila, and F. Iachello, *Phys. Rev. C* **91**, 034304 (2015).
- [26] N. L. Vaquero, T. R. Rodríguez, and J. L. Egido, *Phys. Rev. Lett.* **111**, 142501 (2013).
- EDF [27] J. M. Yao, L. S. Song, K. Hagino, P. Ring, and J. Meng, *Phys. Rev. C* **91**, 024316 (2015).

NME can be decomposed into individual contributions with spin/parity ( $J^\pi$ ) of two decaying neutrons :

$$M^{0\nu} = \sum_J \langle 0_f^+ | \sum_{i \leq j, k \leq l} M_{ij,kl}^J [(\hat{a}_i^\dagger \hat{a}_j^\dagger)^J (\hat{a}_k \hat{a}_l)^J]^0 | 0_i^+ \rangle$$



What can be physical implications ?

What can be foreseen ?

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1. Neutrinoless double beta decay  $^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$
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The model space is truncated to the  $f_{7/2}$  only, for the purpose of schematic studies of certain basic features.

$^{48}\text{Ca}$  ground state  $\Rightarrow f_{7/2}$  closed shell

This is also the condensate of  $J=0^+$  pair (4 pairs).

We can take  $f_{7/2}$  part of the dbd operator.

$$M^{0\nu} = \sum_J \langle 0_f^+ | \sum_{i \leq j, k \leq l} M_{ij,kl}^J [(\hat{a}_i^\dagger \hat{a}_j^\dagger)^J (\hat{a}_k \hat{a}_l)^J]^0 | 0_i^+ \rangle$$

$$i, j, k, l = f_{7/2}$$



## $f_{7/2}$ single orbit model

$^{48}\text{Ca}$  ground state :  $(S^+)^4 |0\rangle$  closed shell

$S^+_{n(p)}$  :  $J=0$  pair of neutrons (protons)

$D^+_{n(p)}$  :  $J=2$  pair of neutrons (protons)

Neutrinoless double beta decay operator  $x_0 S^+_p S_n + x_2 D^+_p D_n + \dots$

$$S^+_p S_n (S^+_n)^4 |0\rangle \sim + S^+_p (S^+_n)^3 |0\rangle$$

$$D^+_p D_n (S^+_n)^4 |0\rangle \sim - D^+_p D_n (S^+_n)^2 |0\rangle$$

$x_0 = -3.32, x_2 = -2.37$   
 same sign

because  $D_n (S^+_n)^4 |0\rangle \sim - D^+_n (S^+_n)^2 |0\rangle$

opposite sign

Quadrupole-quadrupole proton-neutron interaction gives  $^{48}\text{Ti}$  ground state as

$$(c_0 S^+_p S^+_n + c_2 D^+_p D^+_n) (S^+_n)^2 |0\rangle$$

with  $c_0$  and  $c_2$  both positive

Thus, there is always cancellation between  $S^+_p S_n$  and  $D^+_p D_n$  contributions.

Monopole Pairing interaction  $V_M$ ,  $T=1$   $J=0$  part of the interaction.

$$\begin{aligned} V_M &= - G_0 \left( \frac{1}{2} \right) ([a_j^+ a_j^+]^{(0)} [a_{\sim j} a_{\sim j}]^{(0)}) \\ &= - G_0 (S^+ S_{\sim}) \end{aligned}$$

where  $G_0 = - \langle j^2; J=0 | V | j^2; J=0 \rangle$  usually  $G_0 > 0$  (2.24 for GXPF1B)

Quadrupole Pairing interaction  $V_Q$ ,  $T=1$   $J=2$  part of the interaction.

$$\begin{aligned} V_Q &= - G_2 \left( \frac{1}{2} \right) ([a_j^+ a_j^+]^{(2)} [a_{\sim j} a_{\sim j}]^{(2)}) \\ &= - G_2 (D^+ D_{\sim}) \end{aligned}$$

where  $G_2 = - \langle j^2; J=2 | V | j^2; J=2 \rangle$  usually  $G_2 > 0$  (0.94 for GXPF1B)

We can verify some properties,  
by tuning the strength of the quadrupole pairing (Q-pairing)  
interaction  $V_Q$ .

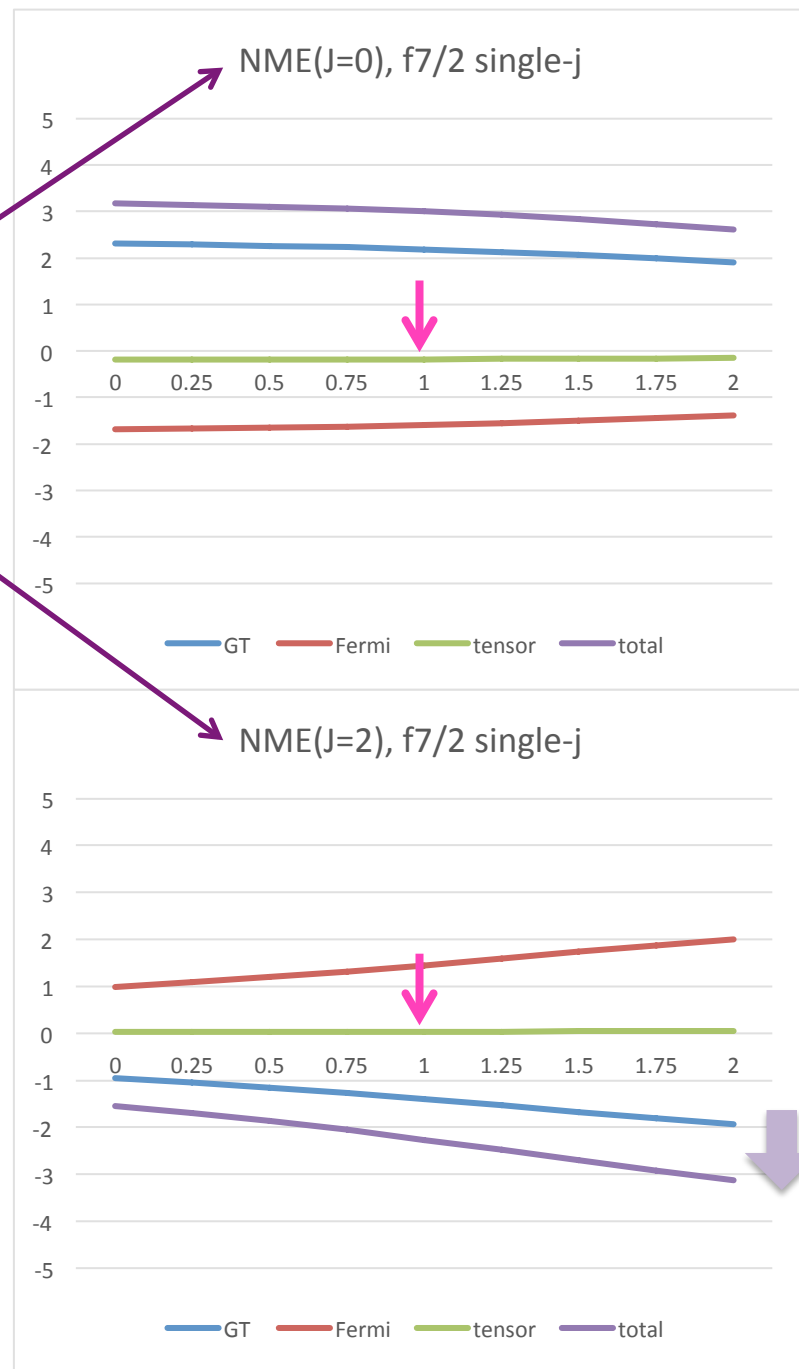
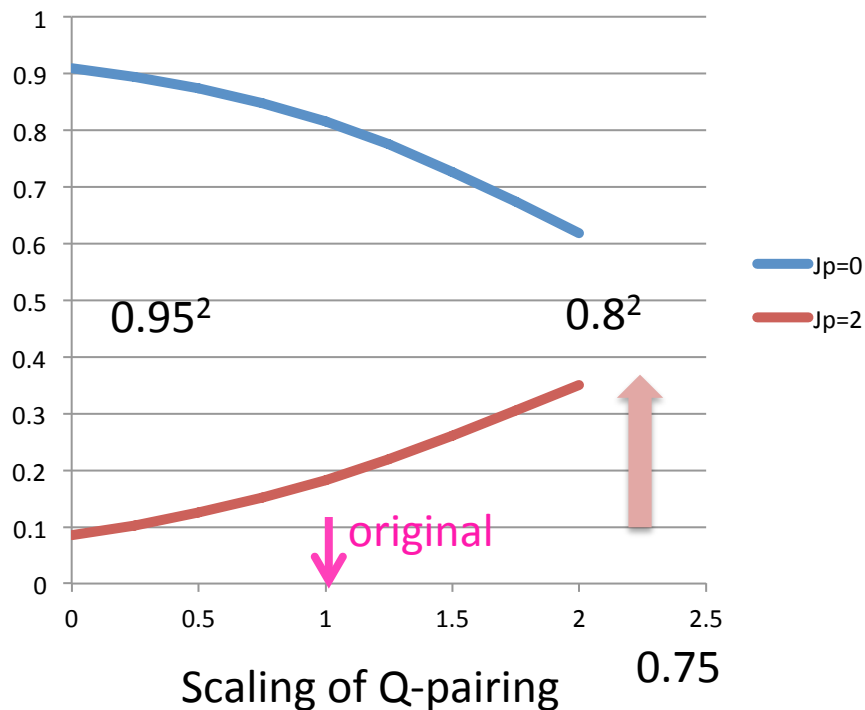
The ground state of  $^{48}\text{Ti}$

$$= \{ (c_0 S_p^+ S_n^+ + c_2 D_p^+ D_n^+) (S_n^+)^2 + \dots \} |0\rangle$$

# of the proton pair = this probability

➡ change of NME by J=0 and 2

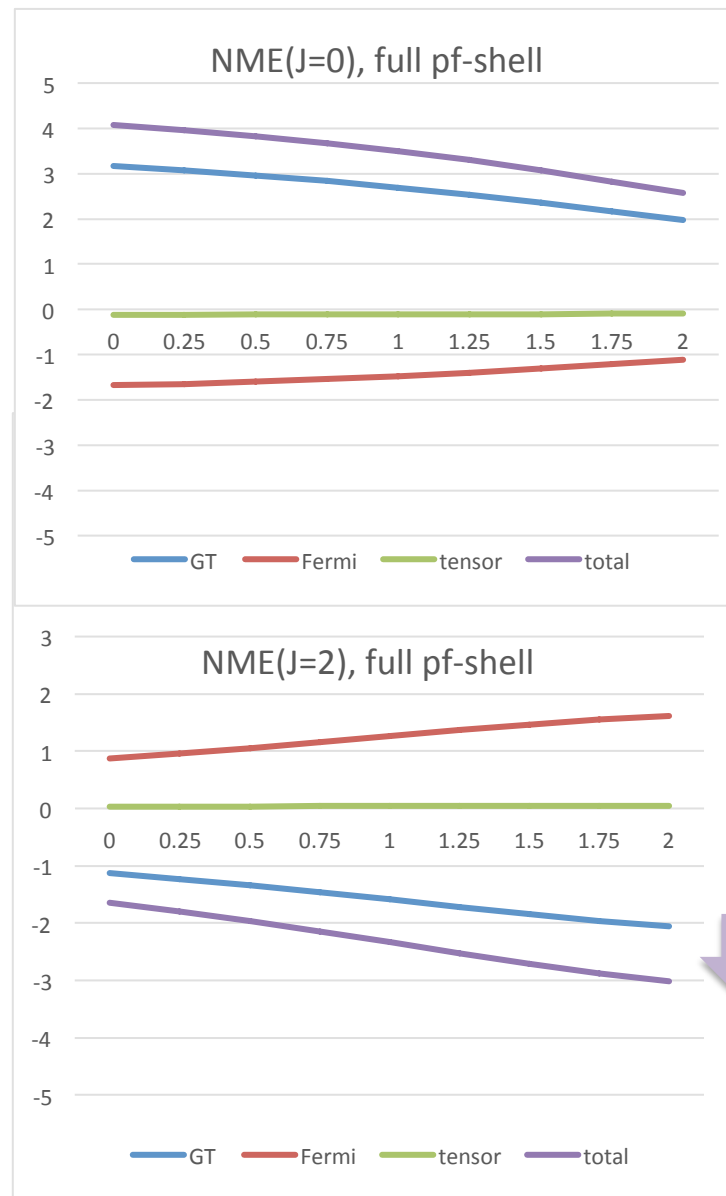
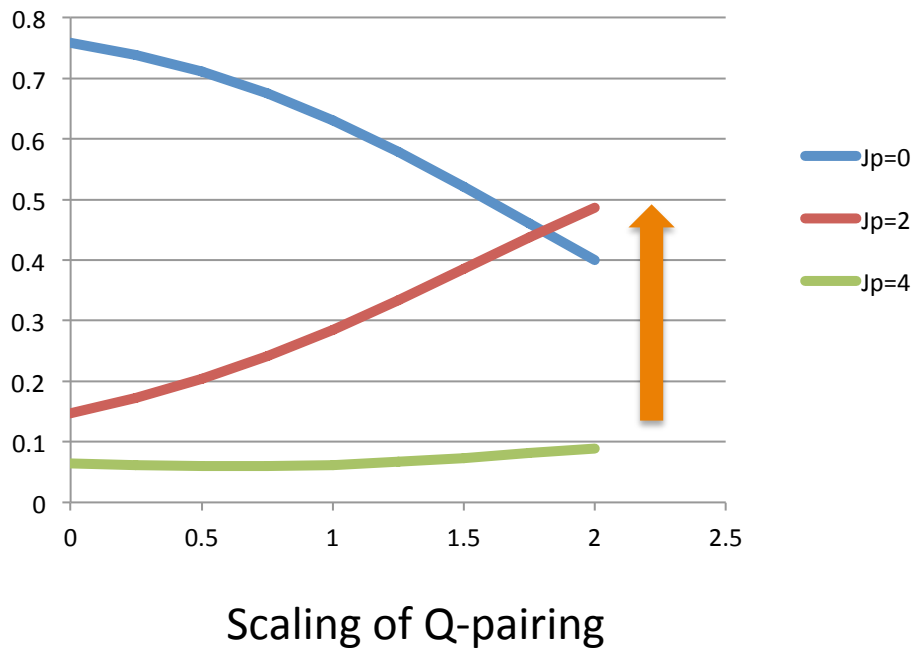
proton pair  $^{48}\text{Ti}$  : f7/2 only



This feature can be seen in the full pf shell calculation

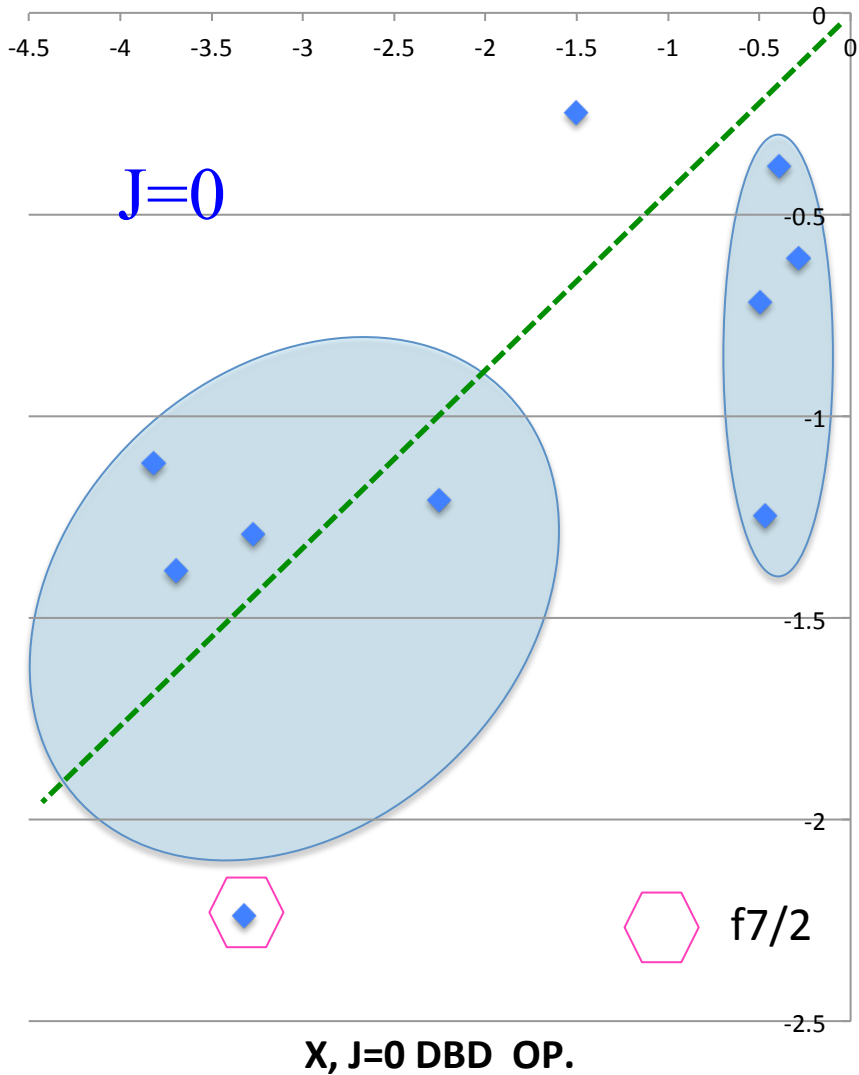
$$\left. \begin{aligned}
 S \text{ pair} &= \sum_i \alpha_i S_i \\
 D \text{ pair} &= \sum_{i < j} \beta_{ij} D_{ij}
 \end{aligned} \right\} \text{coherent sum}$$

The  $J=2$   $T=1$  TBME's are scaled.  
 proton pair  $^{48}\text{Ti}$  : pf shell

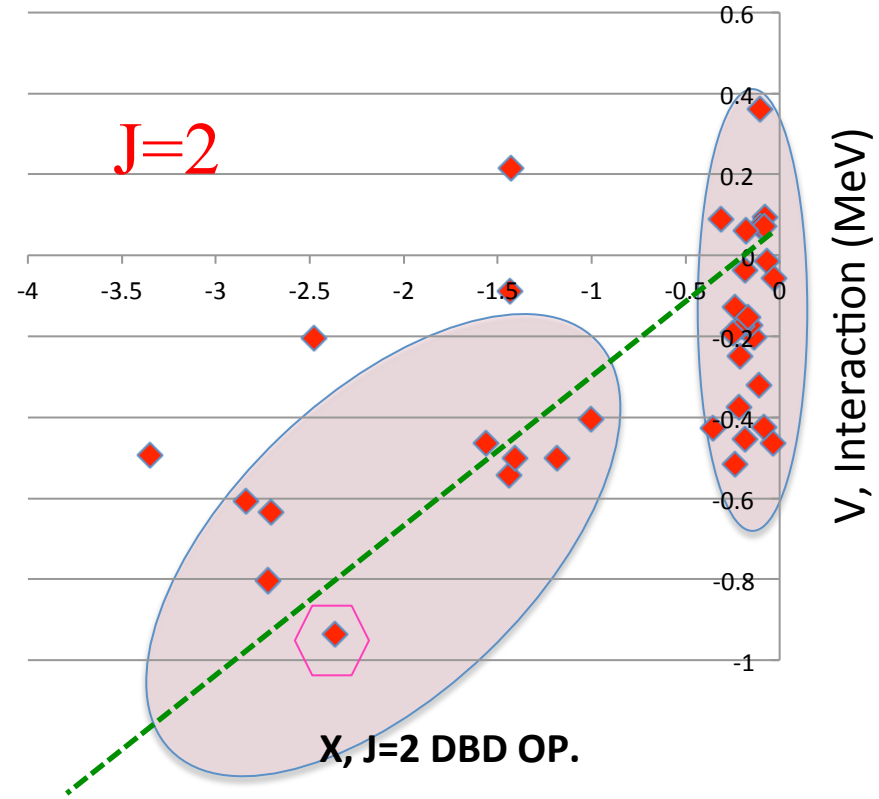


# Correlation between TBME's of Interaction and double-beta decay

$$\langle i j ; J | V \text{ or } X | k l ; J \rangle$$



V, Interaction (MeV)



Two categories  
 (a) correlations  
 (b) small values for dbd

Signs are basically the same

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## Spherical ground state in open shell

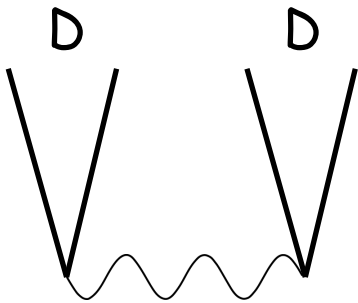
ground state :  $(S^+)^k |0\rangle$       condensation of  $J=0$  pairs

Quadrupole-pairing neutron-neutron interaction mixes as

$$c_1 (S^+_n)^k - c_2 (D^+_n D^+_n) (S^+_n)^{k-2} |0\rangle$$

with  $c_1$  and  $c_2$  both positive.

Ground state correlation (backward scattering) amplitude.



amplitudes with opposite sign between  
neutron-neutron Q-pairing interaction  
and  
neutron-neutron quadrupole-quadrupole  
interaction (which is much weaker)

This is a robust property determined by the double commutator.

**Quadrupole-pairing neutron-neutron** interaction mixes as

$$c_1 (S_n^+)^k - c_2 (D_n^+ D_n^+) (S_n^+)^{k-2} |0\rangle$$

with  $c_1$  and  $c_2$  both positive.

Double-beta decay from this state tends to lead to

$$x_1' S_p^+ (S_n^+)^{k-1} - x_2' (D_p^+ D_n^+) (S_n^+)^{k-2} |0\rangle$$

with  $x_1'$  and  $x_2'$  both positive

**Another source of cancellation**

**Quadrupole-quadrupole proton-neutron** interaction mixes as

$$(c_0 S_p^+ S_n^+ + c_2 D_p^+ D_n^+) (S_n^+)^2 |0\rangle$$

with  $c_0$  and  $c_2$  both positive



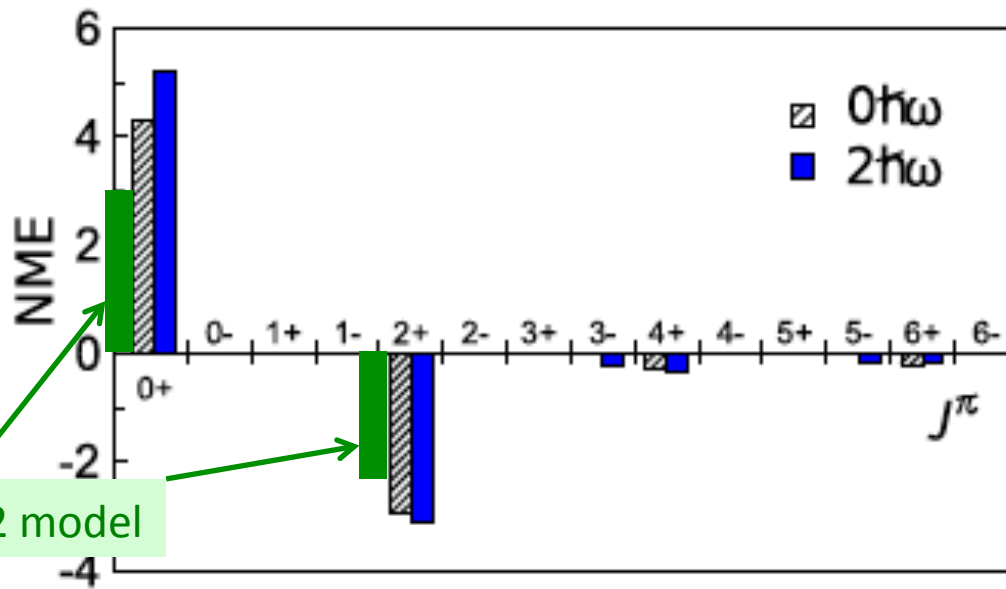
assuming Q-pairing  $J=2$  pair  $\sim [Q, S^+] \sim J=2$  dbd pair

M-pairing  $J=0$  pair  $\sim J=0$  dbd pair



Contributions to NME increase in magnitude :

$$M^{0\nu} = \sum_J \langle 0_f^+ | \sum_{i \leq j, k \leq l} M_{ij,kl}^J [(\hat{a}_i^\dagger \hat{a}_j^\dagger)^J (\hat{a}_k \hat{a}_l)^J]^0 | 0_i^+ \rangle$$



single f 7/2 model

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# Rotational nuclei

Intrinsic state of the ground-state rotational band

Particle-number projected HFB state

$$(\Lambda_p^+)^{K_p} (\Lambda_n^+)^{K_n} |0\rangle$$



Double beta decay

$$(\Lambda_p^+)^{K_{p+1}} (\Lambda_n^+)^{K_{n-1}} |0\rangle$$

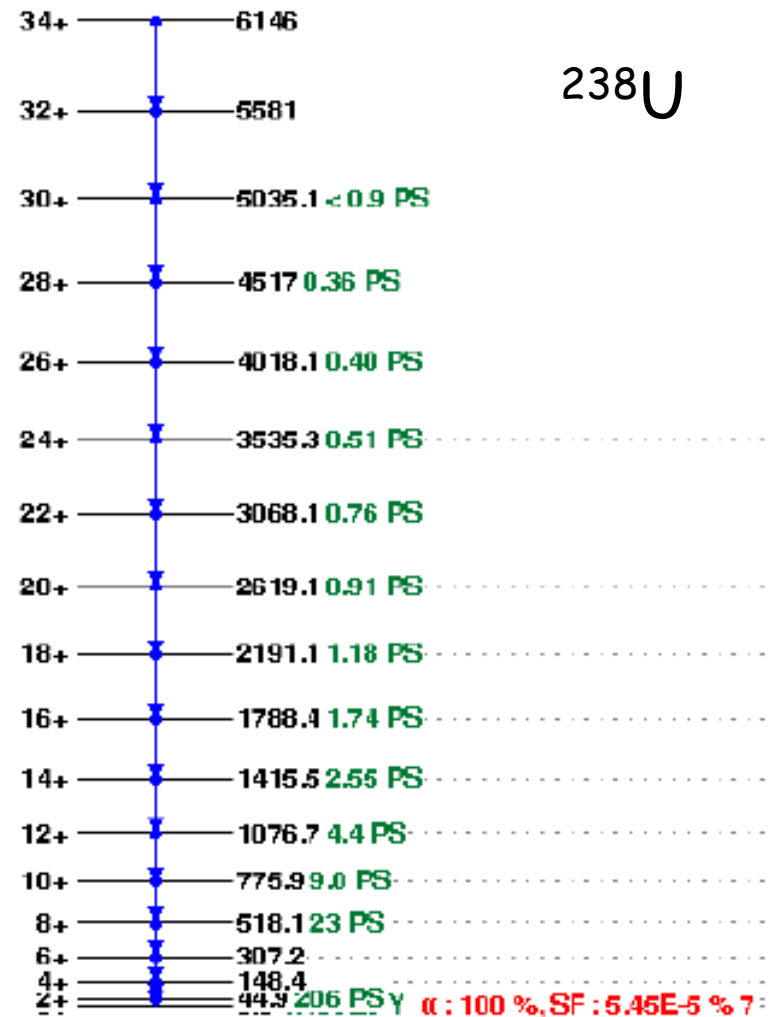
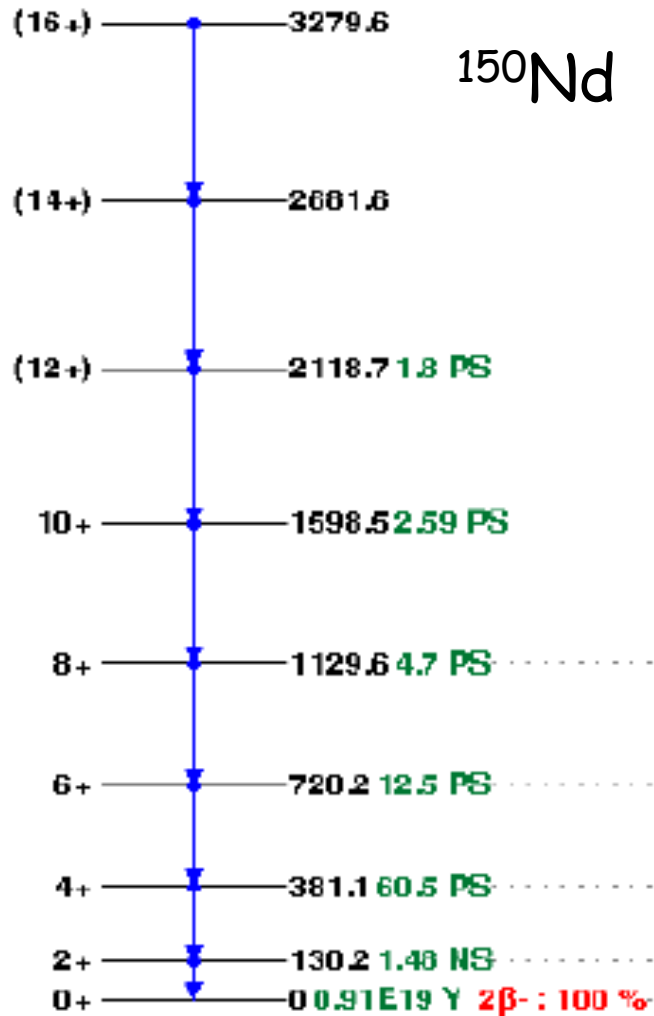
deformed pair :  $\Lambda_{p,n}^+ = \xi_0 S_{p,n}^+ + \xi_2 D_{p,n}^+ + \xi_4 G_{p,n}^+ + \dots$

Higher order terms may contribute, but HFB with  $N$  and  $J^P$  projections will include all effects.

Possibility of a larger NME ... good for establishing  $0\nu$  dbd

# Candidates $^{150}\text{Nd}$ , $^{238}\text{U}$

From the standard list  $^{48}\text{Ca}$ ,  $^{76}\text{Ge}$ ,  $^{82}\text{Se}$ ,  $^{96}\text{Zr}$ ,  $^{100}\text{Mo}$ ,  $^{116}\text{Cd}$ ,  $^{128}\text{Te}$ ,  $^{130}\text{Te}$ ,  $^{150}\text{Nd}$ ,  $^{238}\text{U}$



END