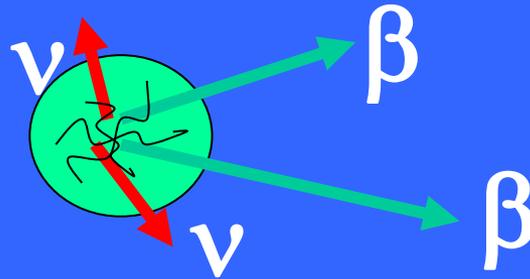
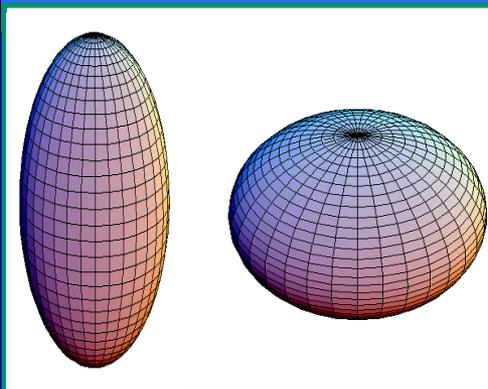


Putting together the pieces of the puzzle in $\beta\beta$ -decay

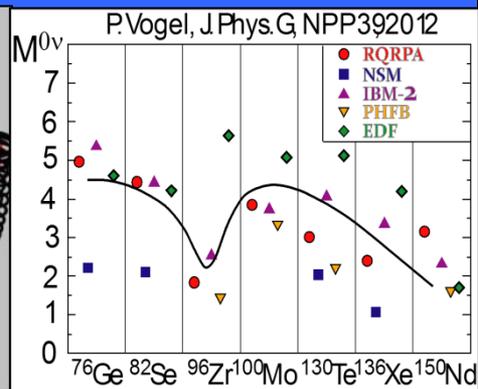
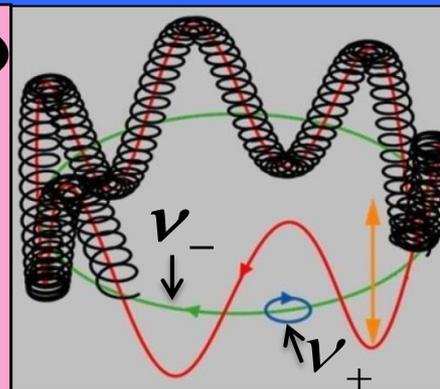


TRIUMF May-2016



Gentle Touch:

- $q_{tr} = 0$
- $\Delta I = 0$
- $0 \hbar\omega$ excitation



The pieces of the puzzle

➤ Chargex-reactions ($^3\text{He},t$) & ($d,^2\text{He}$)

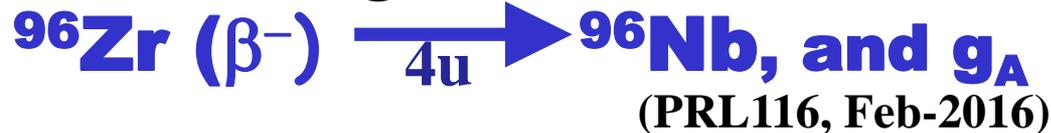
- $2\nu\beta\beta$ nuclear matrix elements
- **Q:** is it relevant for $0\nu\beta\beta$ -decay?
- **A:** yes! (i.e. nucl. shape)

➤ Chargex-reactions

- $0\nu\beta\beta$ nuclear matrix element
- **Q:** is it possible?
- **A:** yes! (NME's & 2^- states, occupation #'s)

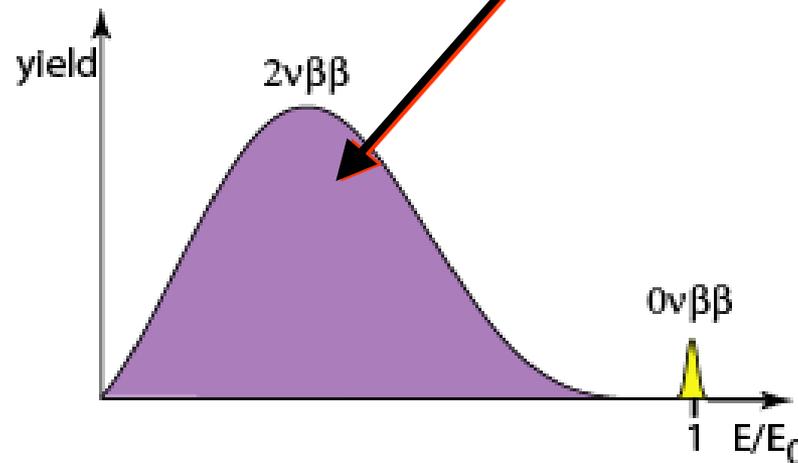
➤ Mass measurements

- $0\nu\beta\beta$ nuclear matrix element
- **Q:** what is the connection?
- **A:** ^{96}Zr is a „golden“ case



N_{ucl.} **M**_{atrix} **E**_{lements}

$2\nu\beta\beta$ - β -decay



**q-transfer like in ordinary
 β -decay**

($q \sim 0.01 \text{ fm}^{-1} \sim 2 \text{ MeV}/c$)

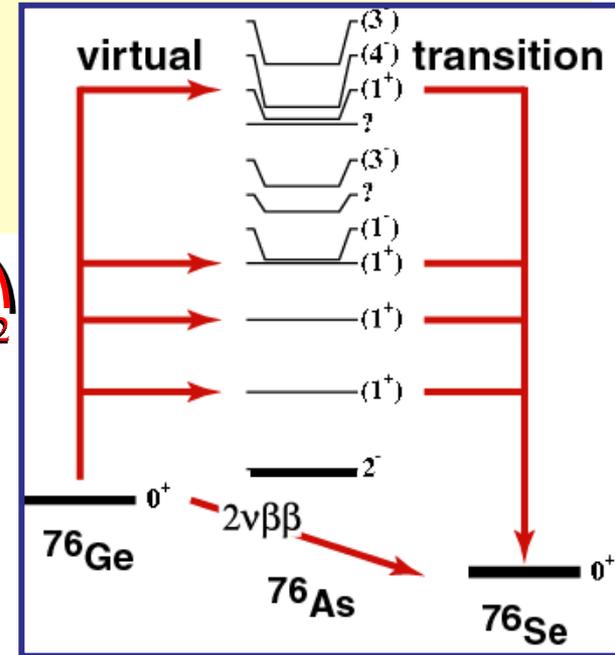
i.e. only allowed transitions possible

$$\Gamma_{(\beta^-\beta^-)}^{2\nu} = \frac{C}{8\pi^7} \left(\frac{G_F g_A}{\sqrt{2}} \cos(\Theta_C) \right)^4 |M_{\text{DGT}}^{(2\nu)}|^2 \mathcal{F}_{(-)}^2 f(\mathbf{Q})$$

$$= G^{2\nu}(\mathbf{Q}, Z) |M_{\text{DGT}}^{(2\nu)}|^2$$

$$\propto Q^{11} \cdot Z^2$$

$\text{exp} \approx 10^{-3} \text{ MeV}^{-2}$
extracted from
half-life

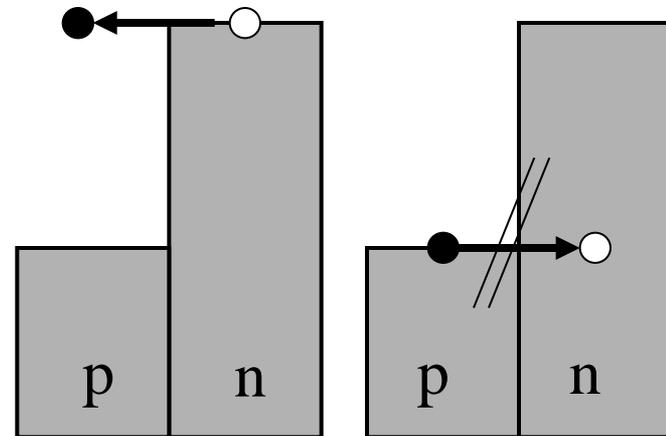


favorable:

1. high Q-value
2. large Z

unfavorable (but cannot be changed):

1. large neutron excess
(Pauli-blocking)



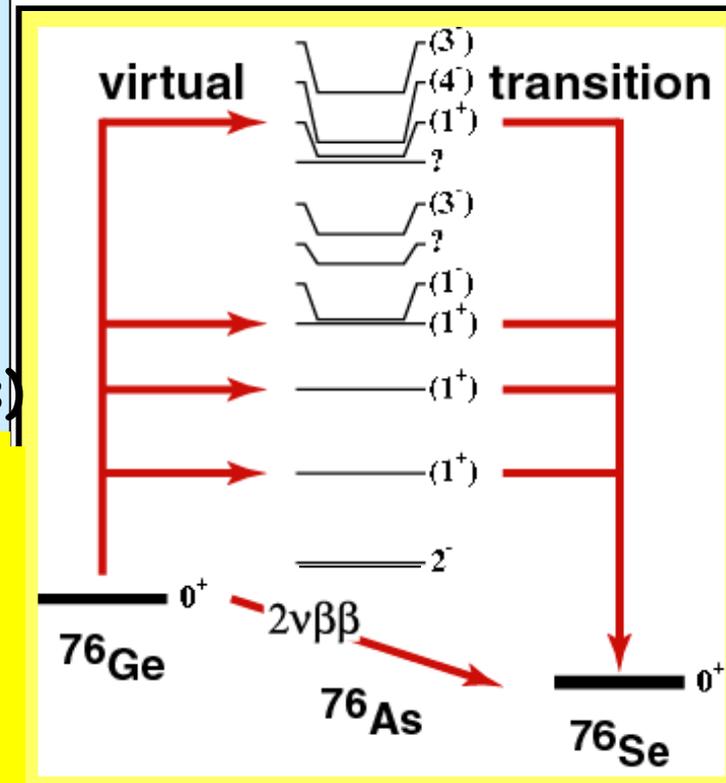
$$M_{\text{DGT}}^{(2\nu)} = \sum_m \frac{\langle \mathbf{0}_{g.s.}^{(f)} | \sum_k \sigma_k \tau_k^- | \mathbf{1}_m^+ \rangle \langle \mathbf{1}_m^+ | \sum_k \sigma_k \tau_k^- | \mathbf{0}_{g.s.}^{(i)} \rangle}{\frac{1}{2} Q_{\beta\beta}(\mathbf{0}_{g.s.}^{(f)}) + E(\mathbf{1}_m^+) - E_0}$$

$$= \sum_m \frac{M_m \quad GT^+ \quad M_m \quad GT^-}{E_m}$$

to remember:

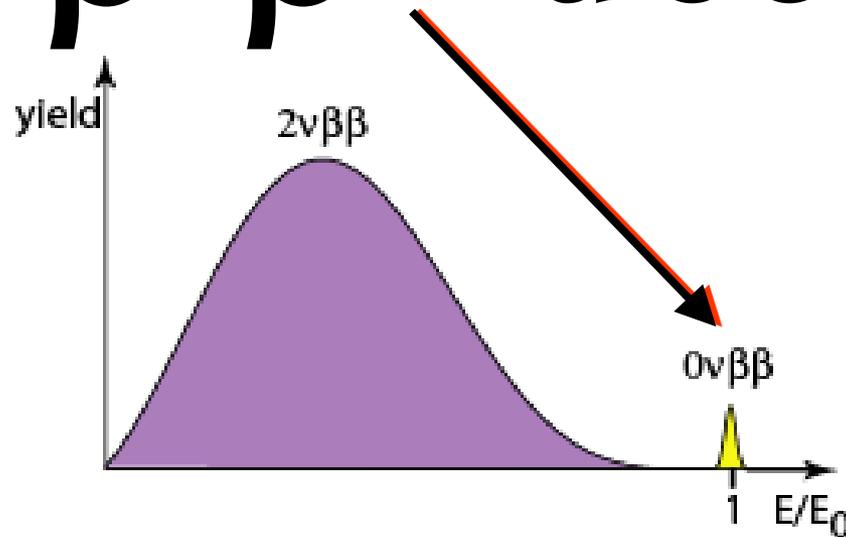
1. 2 sequential & „allowed“ β^- -decays of „Gamow-Teller“ type
2. „1, 2, 3, ... forbidden“ decays negligible
3. Fermi-transitions do not contribute (because of different isospin-multiplets)

Can be determined via charge-exchange reactions in the (n,p) and (p,n) direction (e.g. (d, ^2He) or (^3He , t))



N_{ucl.}**M**_{atrix}**E**_{lements}

$0\nu\beta\beta$ decay



neutrino is a virtual particle

$q \sim 0.5 \text{ fm}^{-1}$ ($\sim 100 \text{ MeV}/c$)

(due to Heisenberg $\Delta q \cdot \Delta x \sim 1$)

degree of forbiddenness is lifted

$$\Gamma_{(\beta^-\beta^-)}^{0\nu} = G^{0\nu}(Q,Z) g_A^4 \left| M_{\text{DGT}}^{(0\nu)} - \left(\frac{g_V}{g_A} \right)^2 M_{\text{DF}}^{(0\nu)} \right|^2 |m_{\nu_e}|^2$$

$$\propto Q^5 \cdot Z^4$$

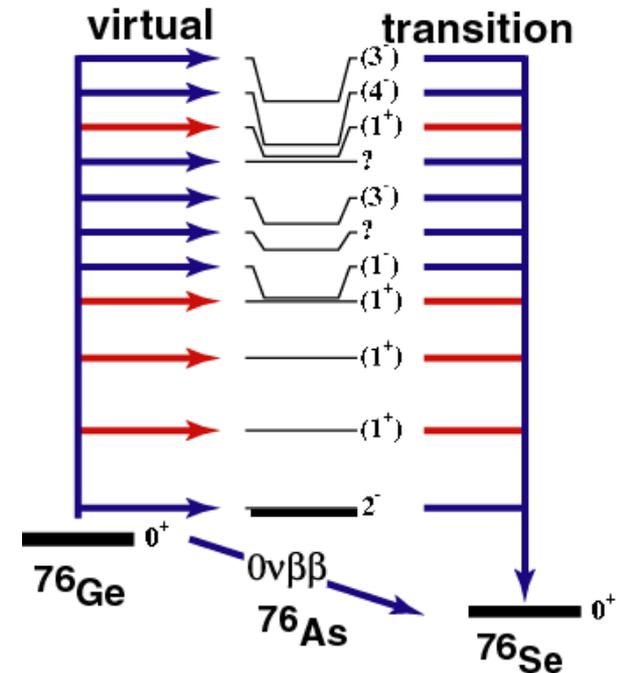
theory
 $\approx 10 !!$
 largely independent of (A,Z)
 (except near magic nuclei)

mass of
 Majorana- ν !

to remember:

1. „higher-fold forbidden“ transitions possible
2. Fermi-transitions important
3. „Pauli-blocking“ largely lifted
4. large Q-value, high Z important

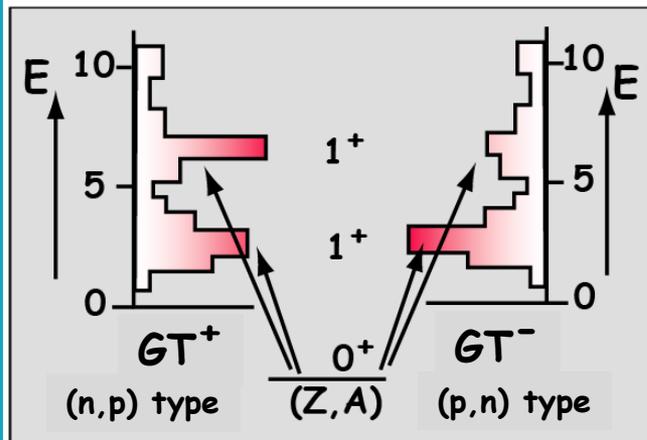
NOT (easily) accessible via charge-exchange reactions



1st piece of puzzle

Charge-exchange reactions

(p,n) type & (n,p) type



Q: what is the connection between „weak $\sigma\tau$ operator“ and the hadronic reaction

A: dominance of the $V_{\sigma\tau}$ effective interaction at medium energies

$$M(GT) = \langle 1^+ || \sigma\tau^\pm || 0_{g.s.}^+ \rangle$$

$$B(GT) = \frac{1}{2J_i+1} |M(GT)|^2$$

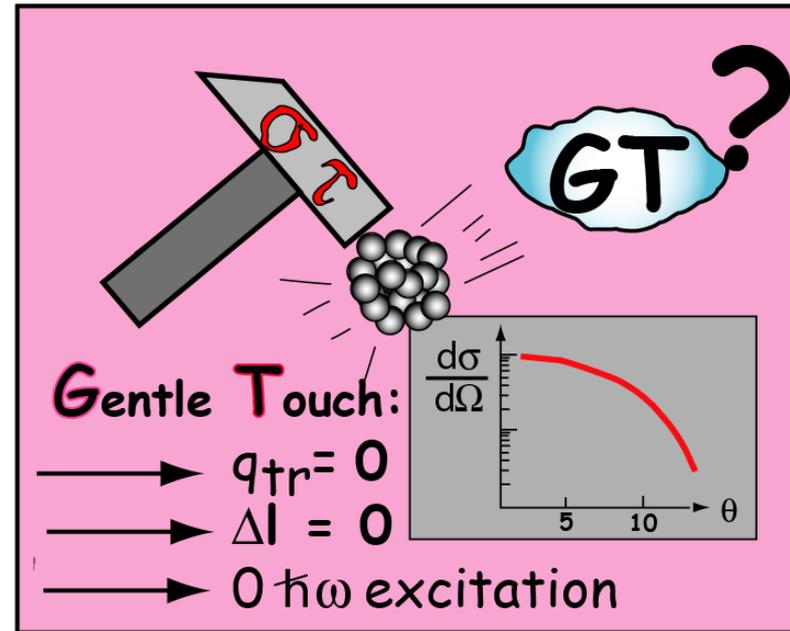
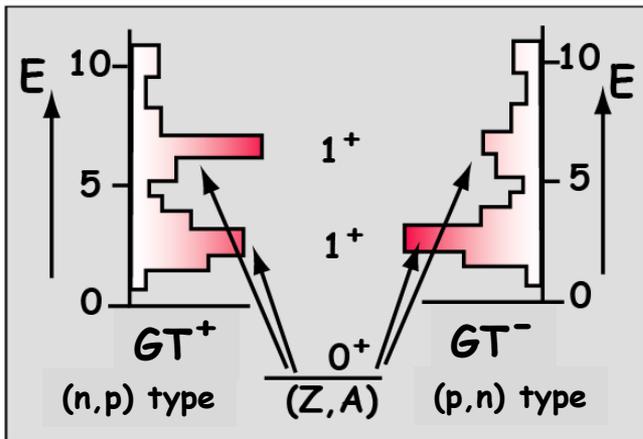
hadronic probes: (n, p), (d, ^2He), (t, ^3He)

or (p, n), (^3He , t)

$$\left[\frac{d\sigma}{d\Omega} \right] = \left[\frac{\mu}{\pi\hbar} \right]^2 \frac{k_f}{k_i} N_d |V_{\sigma\tau}|^2 |\langle f | \sigma\tau | i \rangle|^2$$

$q = 0!!$

largest at 100 - 200 MeV/A



$$M(GT) = \langle 1^+ || \sigma \tau^\pm || 0_{g.s.}^i \rangle$$

$$B(GT) = \frac{1}{2J_i+1} |M(GT)|^2$$

hadronic probes: (n,p), (d,²He), (t,³He)

or (p,n), (³He,t)

$$\left[\frac{d\sigma}{d\Omega} \right] = \left[\frac{\mu}{\pi\hbar} \right]^2 \frac{k_f}{k_i} N_d |v_{\sigma\tau}|^2 |\langle f | \sigma \tau | i \rangle|^2$$

$q = 0!!$

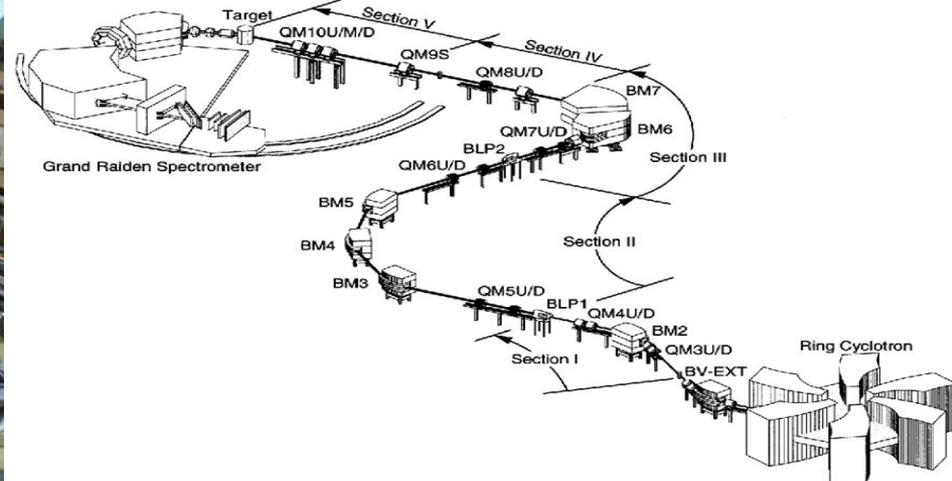
largest at 100 - 200 MeV/A

Charge-exchange reactions

Grand Raiden Magnetic Spectrometer

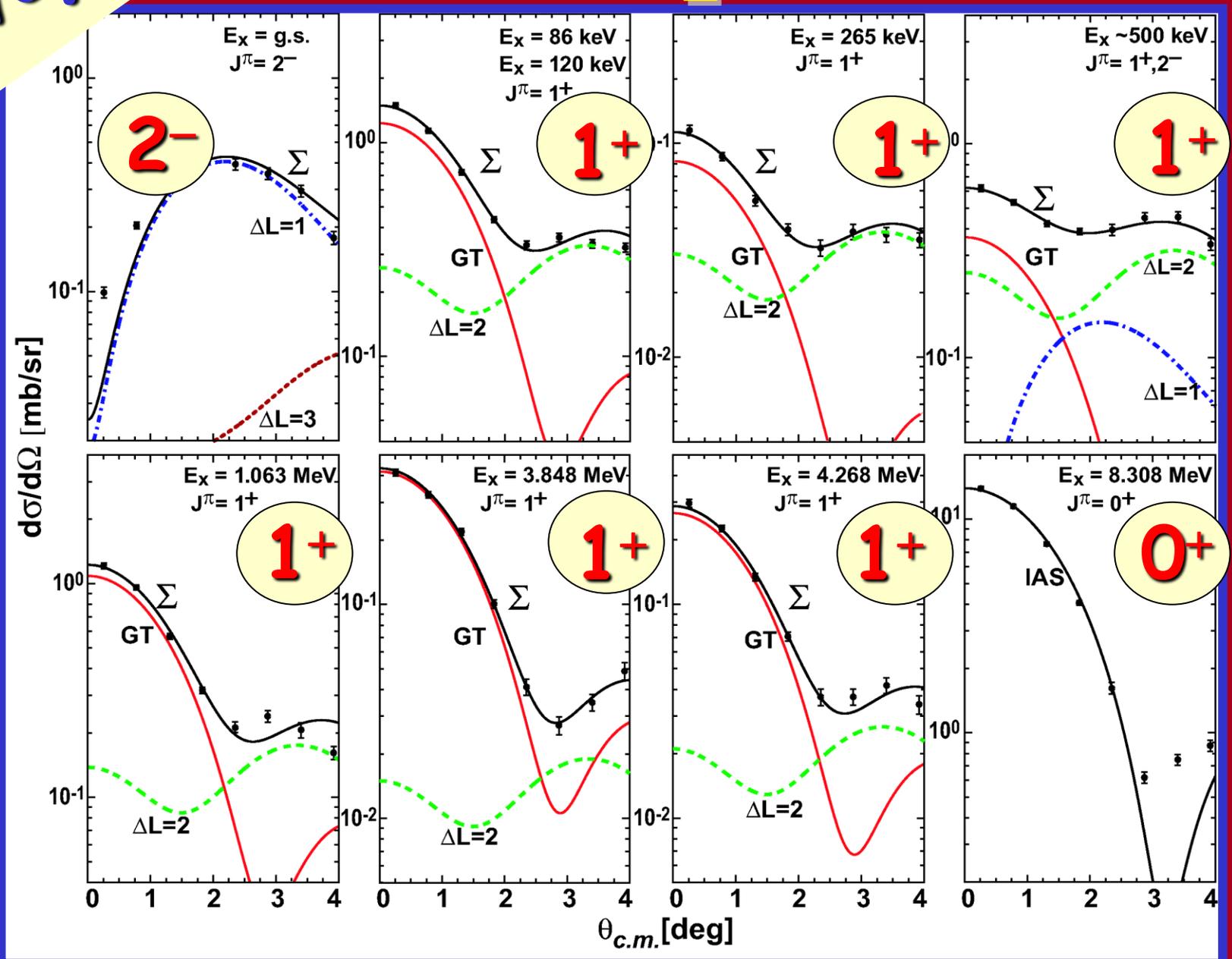


$\Delta E/E \sim 5 \times 10^{-5}$ ~ 25 keV
at 420 MeV (^3He)



$^{76}\text{Ge}-^{76}\text{As}$

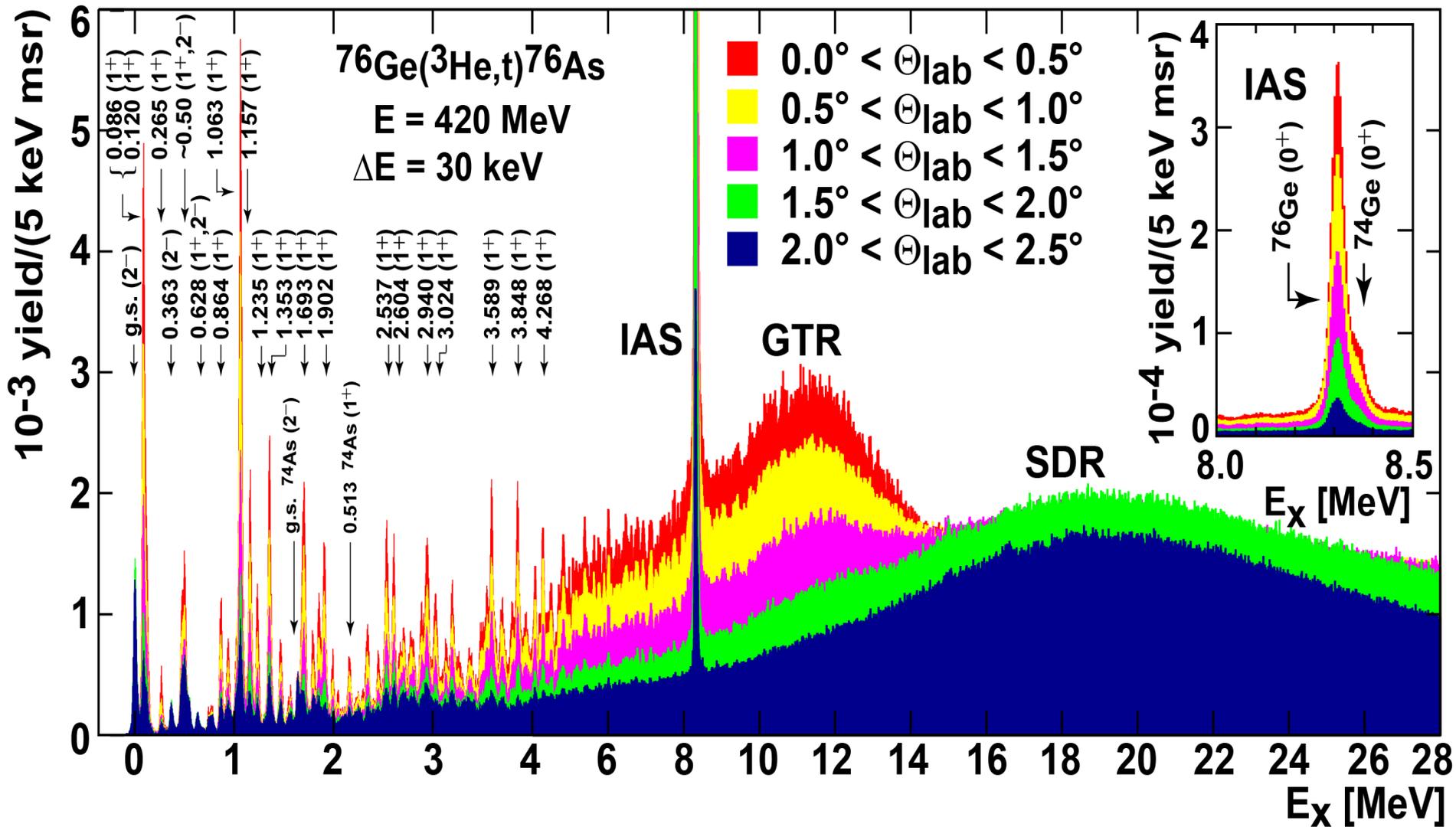
examples



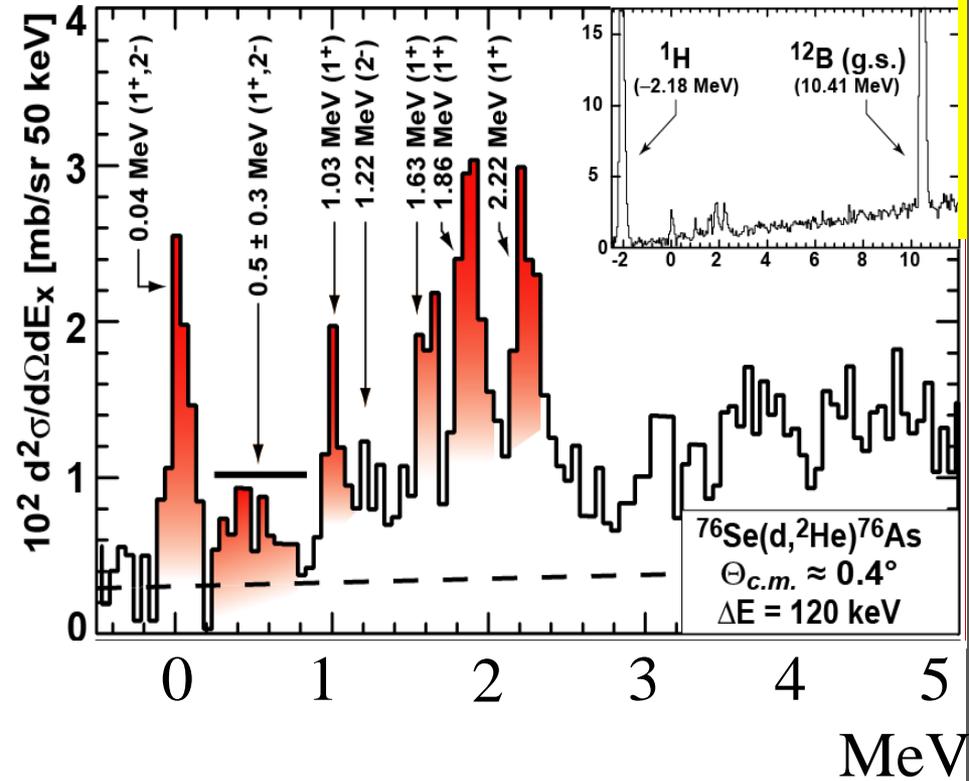
^{76}Ge

$N-Z=10$

Resolution is the key !!!



**almost 70 !! resolved single states up to 5 MeV
 identified as GT 1^+ transitions !!!**



the other leg
 (BGT⁺):
 $^{76}\text{Se}(d, ^2\text{He})^{76}\text{As}$
 ($\Delta E = 120 \text{ keV}$)

Nuclear matrix elements and deformation

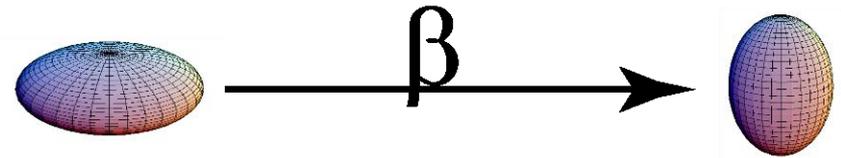
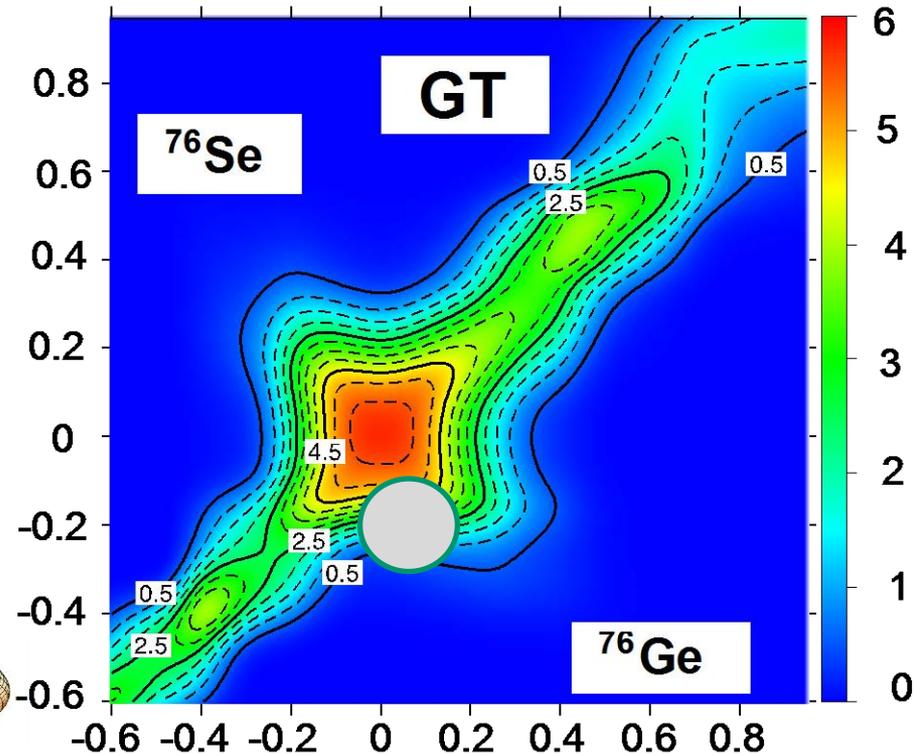
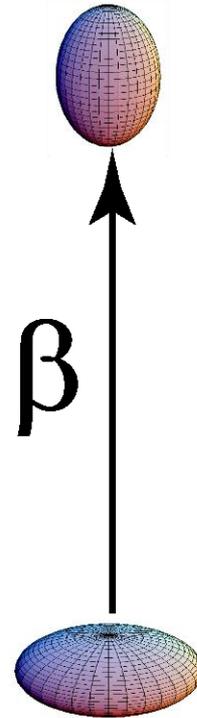
$${}^{76}\text{Ge}: \beta \sim +0.1$$

$${}^{76}\text{Se}: \beta \sim -0.2$$

reduction of the NME
due to deformation is
theoretically confirmed

but

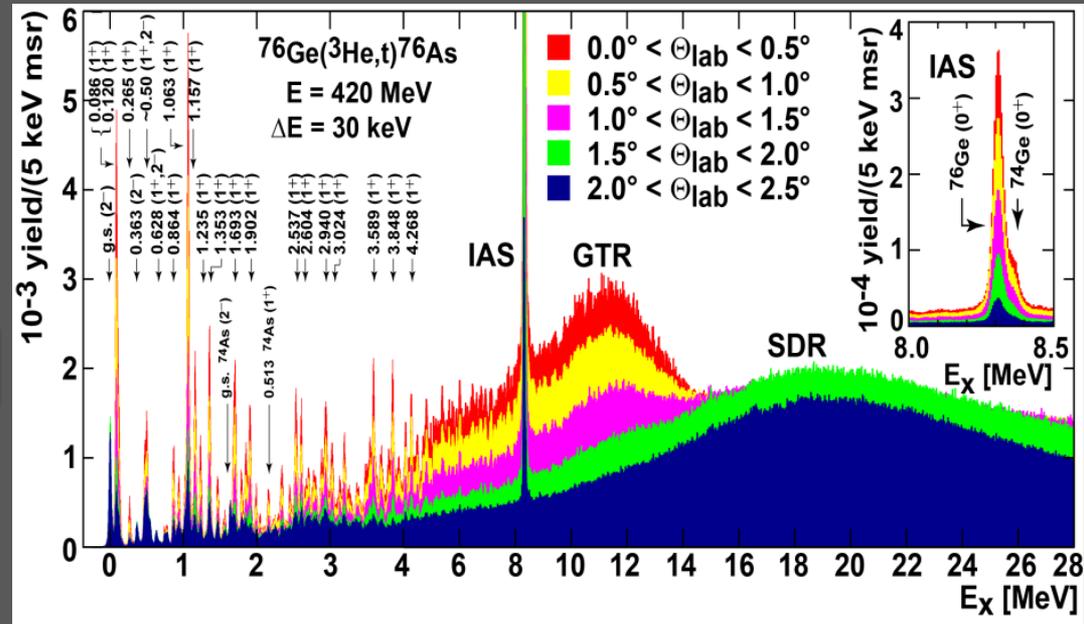
expm'ly it seems to manifests
itself (in $2\nu\beta\beta$ decay) by a lack
of correlation between the two
different $B(\text{GT})$ „legs“, rather
than a reduction of individual
strength



From: T. R. Rodriguez, et al, PRL105 (2010)

another surprise:

low-E part of
NME makes up
~100% of total
 $2\nu\beta\beta$ -ME



no need for
GT giant resonance contribution

100 Mo

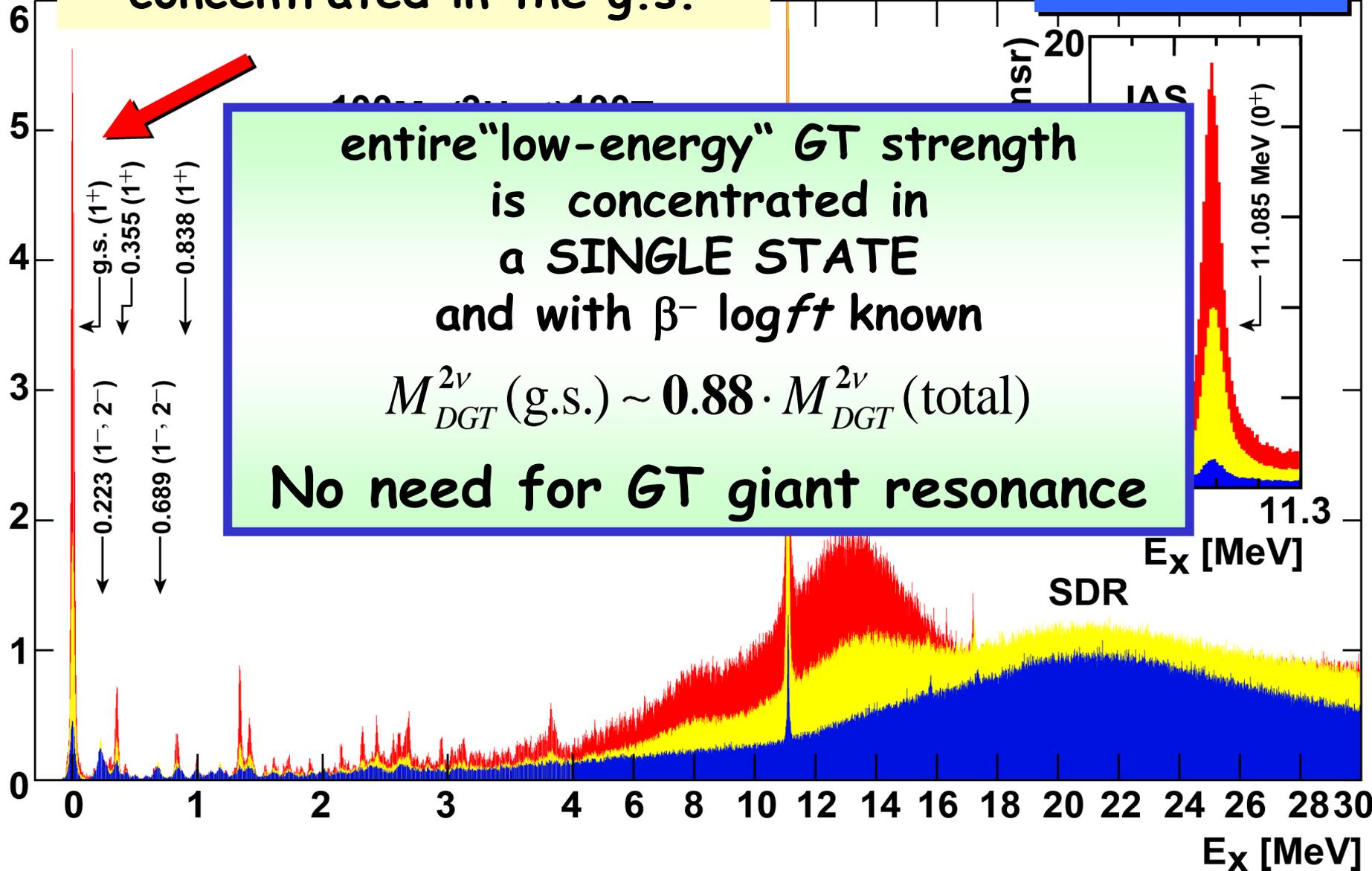
N-Z=16

**also useful as SN neutrino detector
(sensitive to ν temperature in SN)**

HERE: almost the entire low-E GT strength is concentrated in the g.s.

100Mo

10⁻³ yield/(5 keV msr)

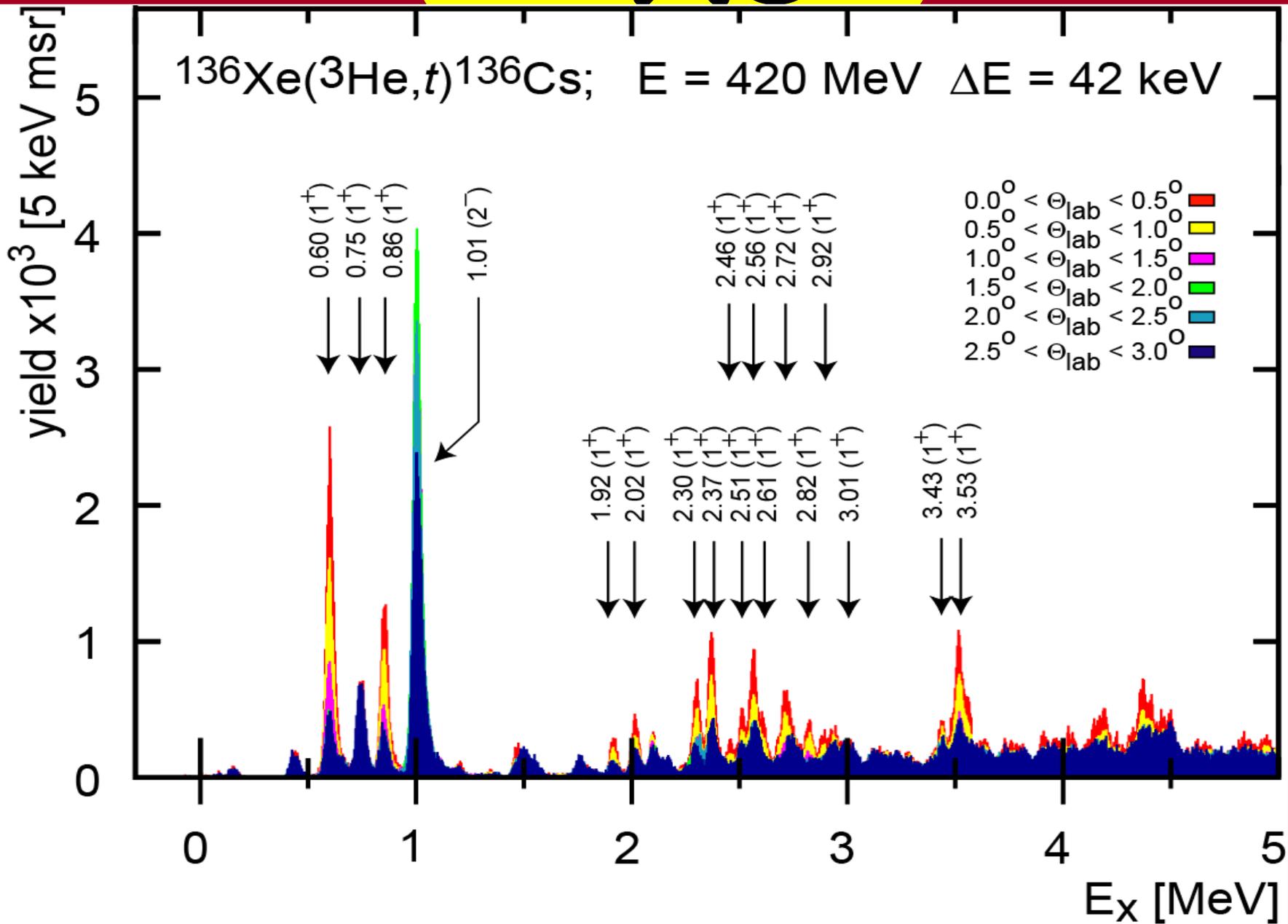


^{136}Xe

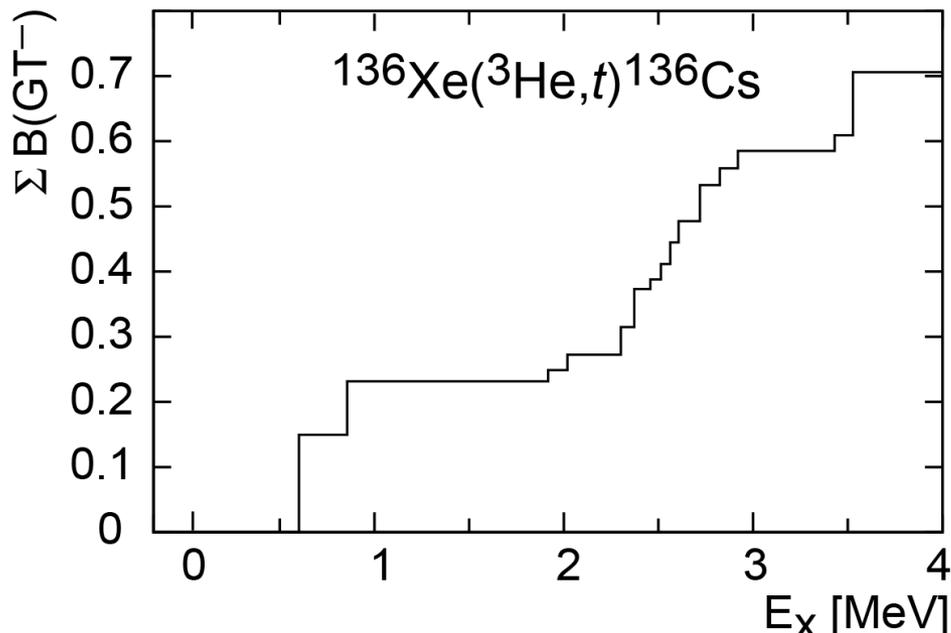
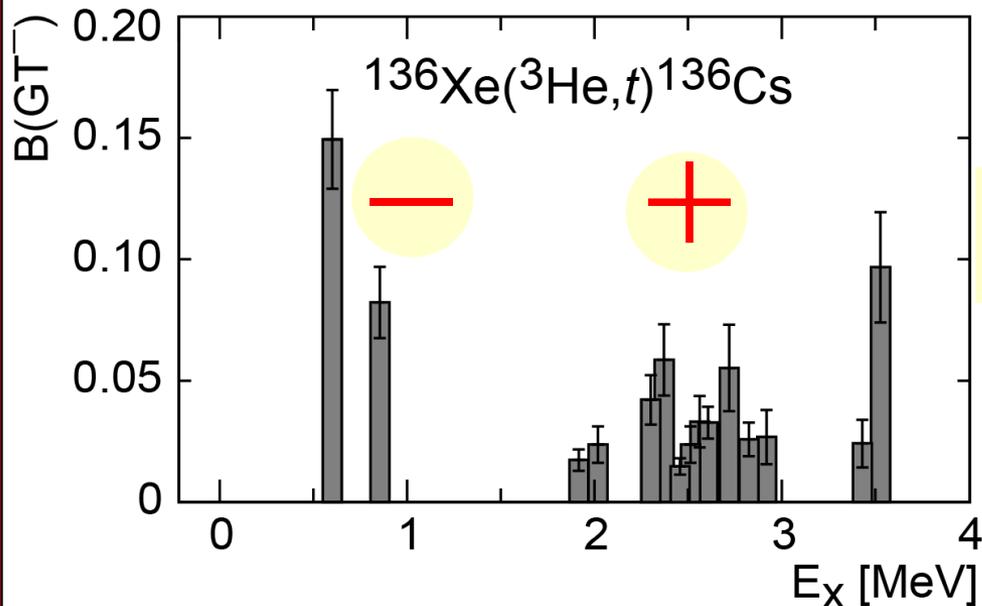
$N-Z=28$

question: why so stable !!!

^{136}Xe



What's the size of the NME?



$$T_{1/2}^{2\nu} = 2.2 \cdot 10^{21} \text{ yr}$$

$$M_{\text{DGT}}^{(2\nu)} \sim 0.019 \text{ MeV}^{-1}$$

all signs positive \rightarrow

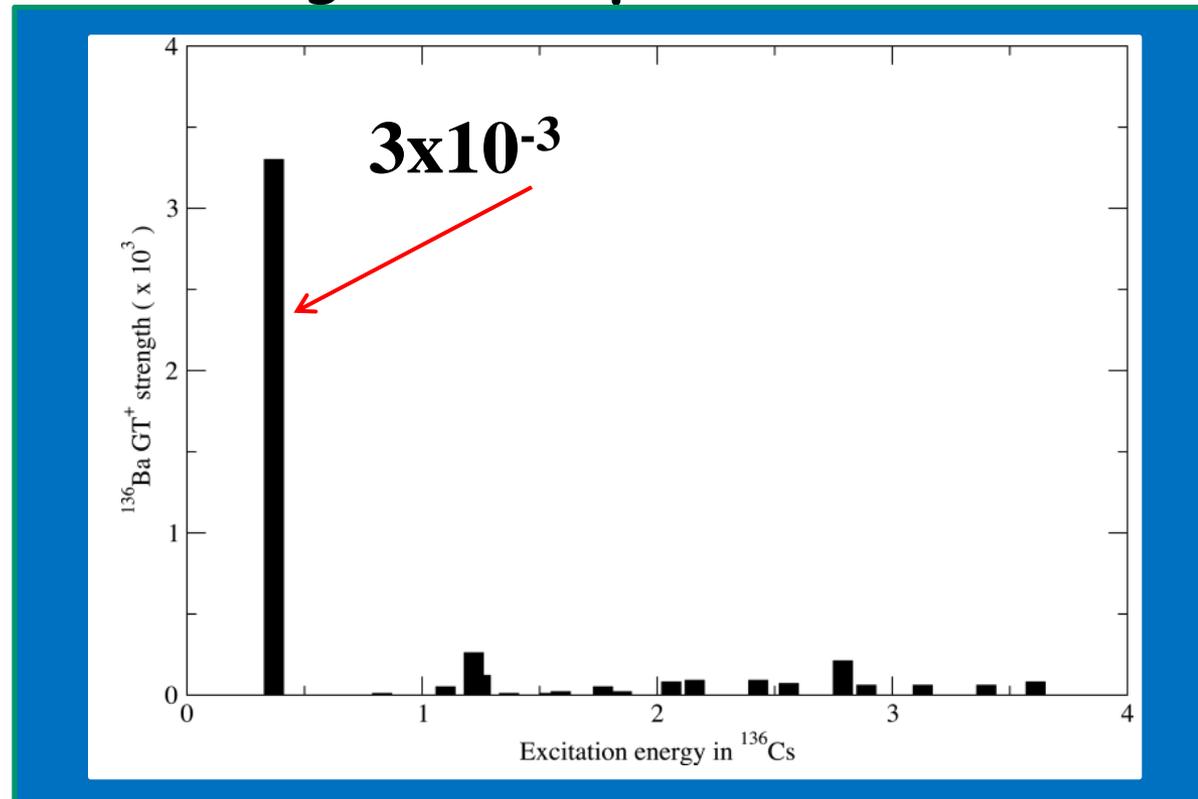
$$B_m \text{ } GT^+ \approx 10^{-2} \cdot B_m \text{ } GT^-$$

$$B_m \text{ } GT^+ \approx 10^{-3} \text{ !!!!}$$

A. Poves (simultaneous to our publication):

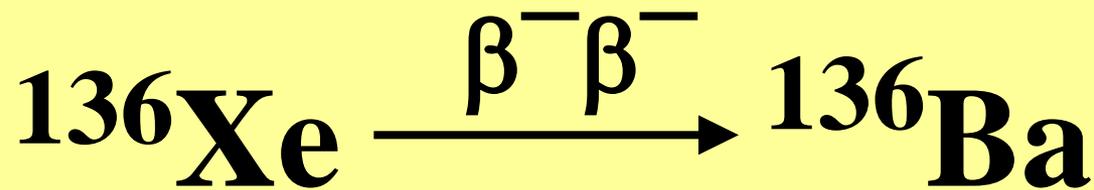
there is no $B(GT^+)$ strength, except for lowest 1^+ state

Recall:
 ^{136}Xe is almost
doubly magic!!



Shell model provides conclusive explanation for the deemed „pathologically“ long half-life of ^{136}Xe .

Expt'l test: $^{136}\text{Ba}(d, ^2\text{He})^{136}\text{Cs}$

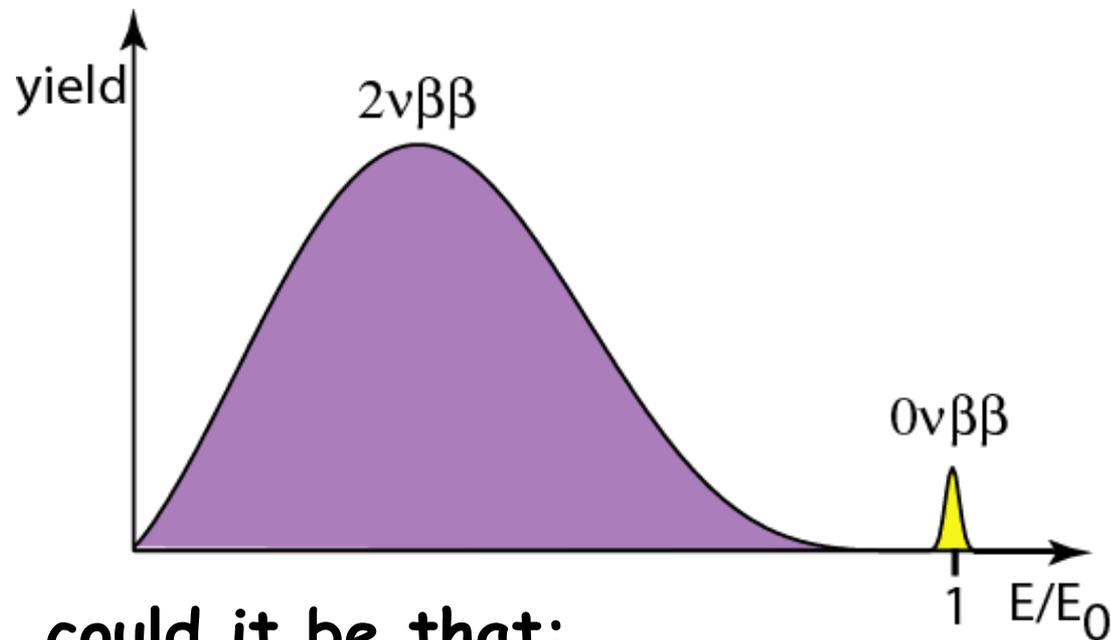


expmt:

$2\nu\beta\beta$ NME is exceptionally small

question:

how does the ME scale in the case of $0\nu\beta\beta$ decay?



could it be that:

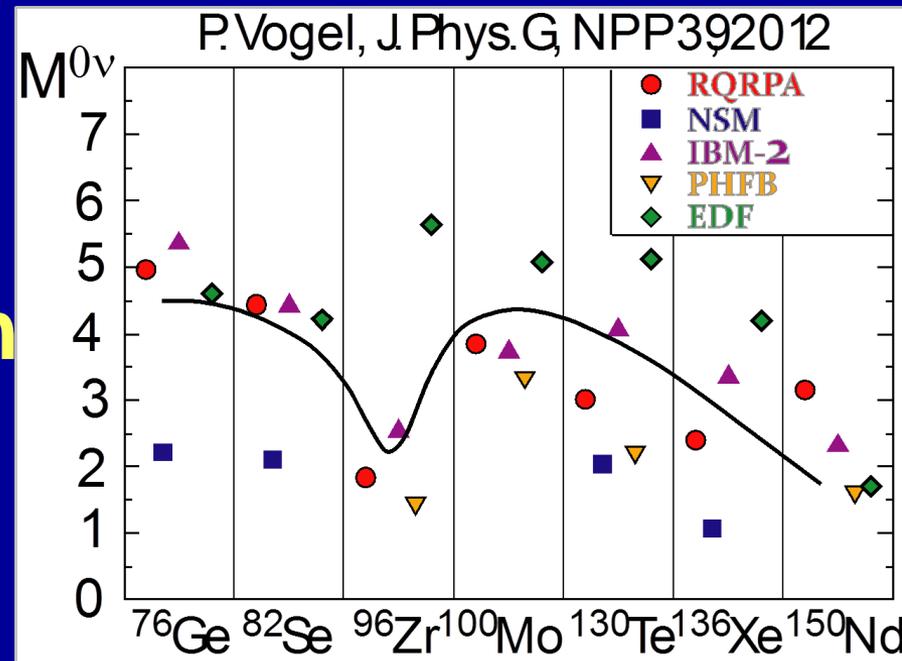
$2\nu\beta\beta$ ME is suppressed **AND**

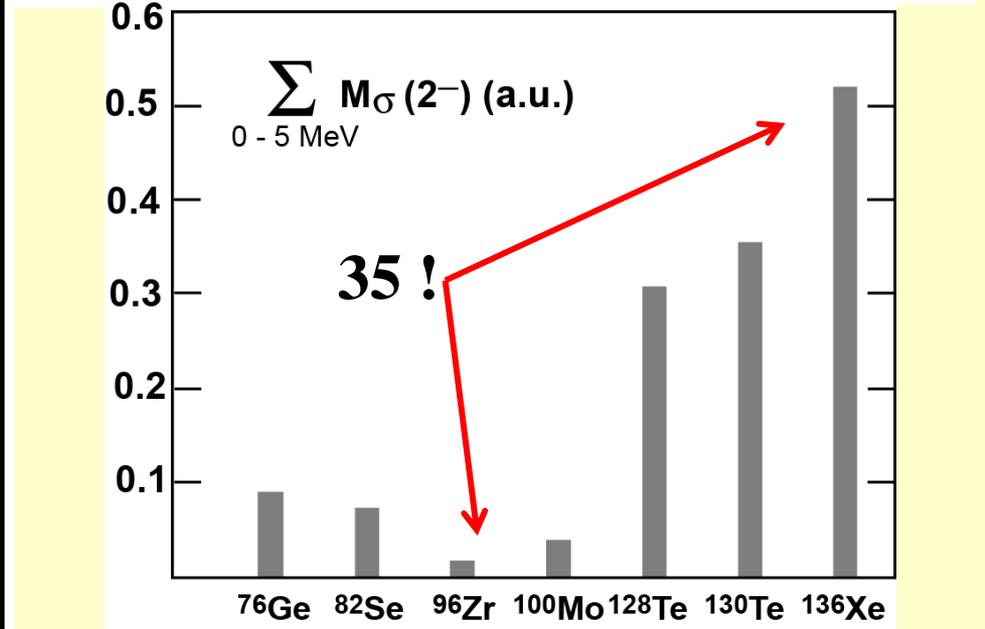
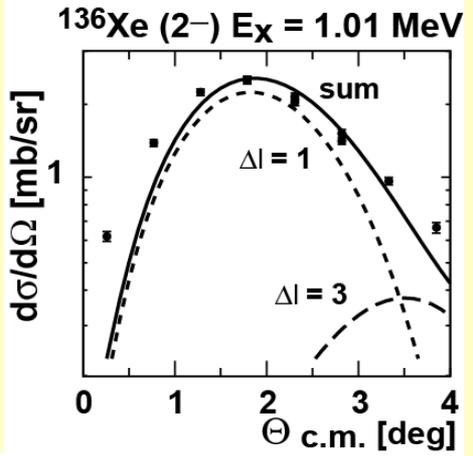
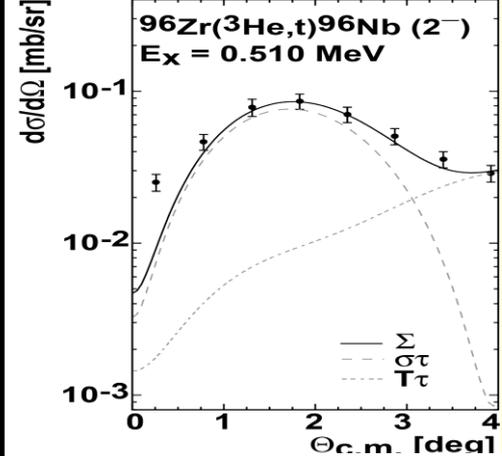
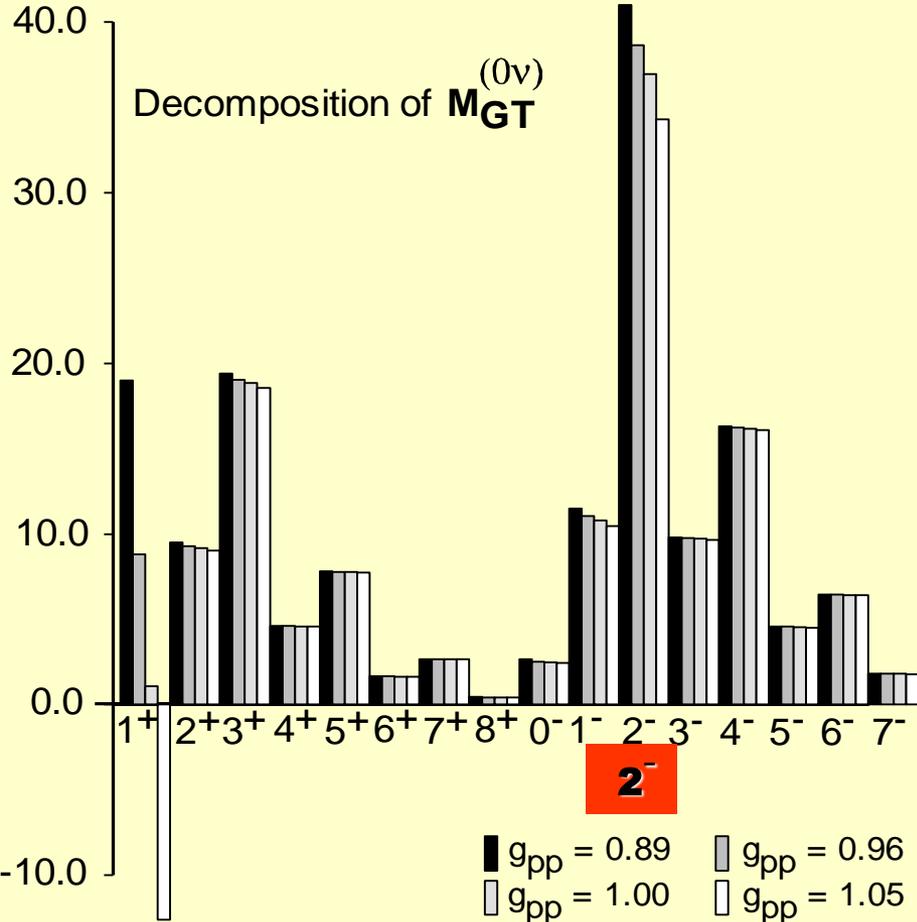
$0\nu\beta\beta$ ME is enhanced ???

2nd piece of puzzle

Charge-exchange reaction towards the $0\nu\beta\beta$ NMEs ???

Here:
2⁻ states and occupation
vacancy numbers
via chargex reactions





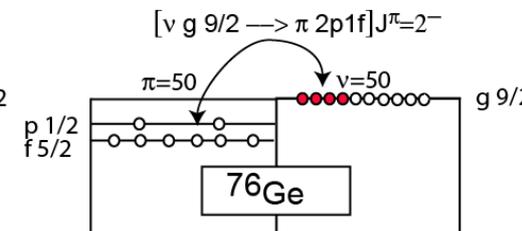
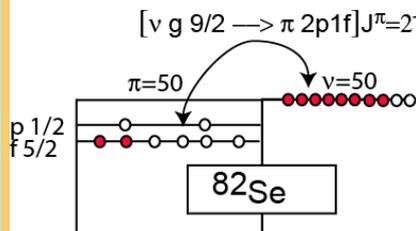
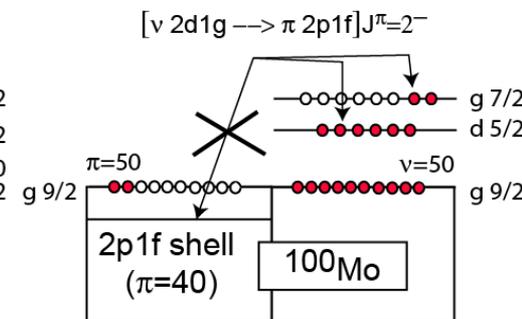
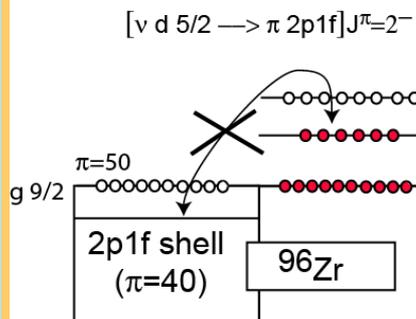
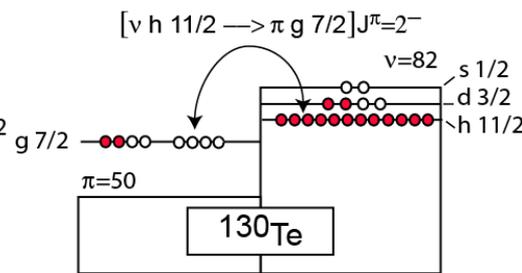
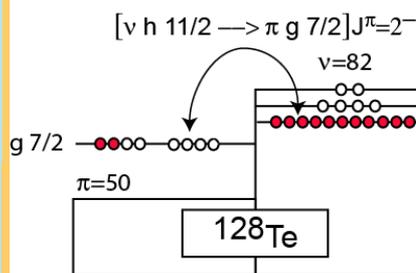
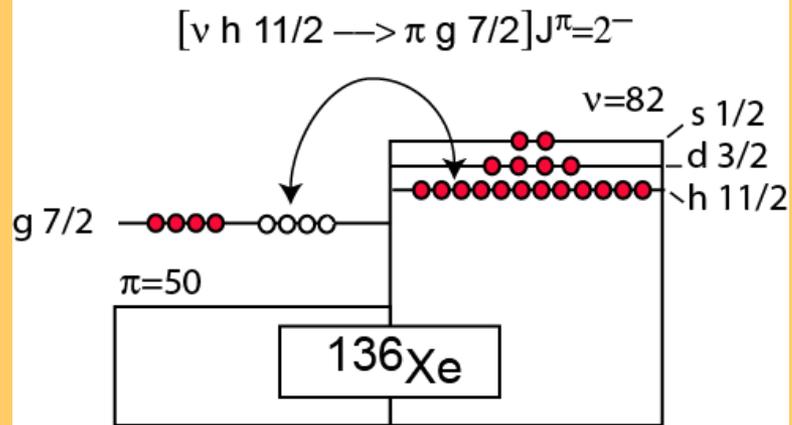
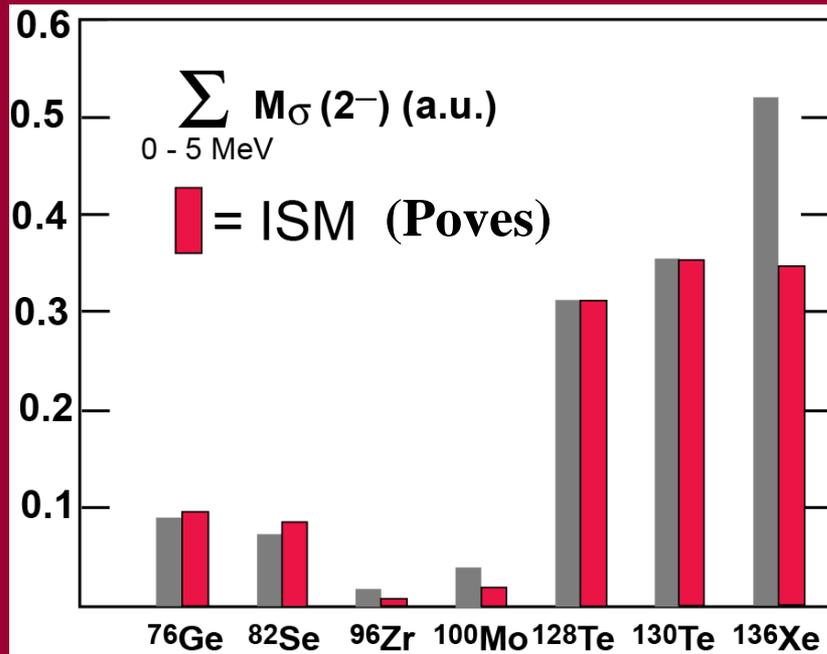
Theory:

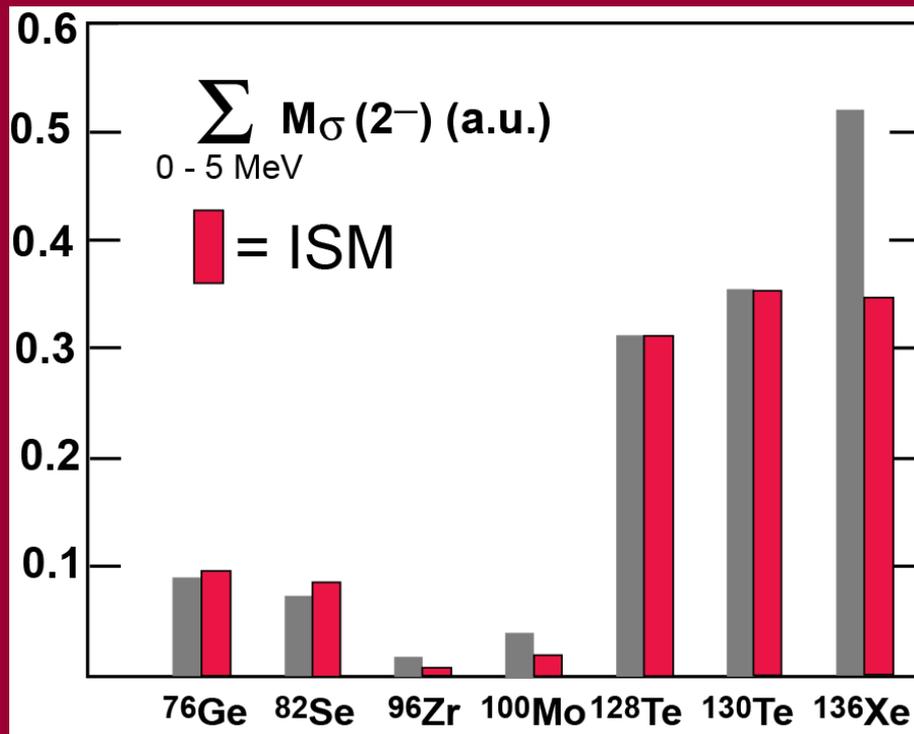
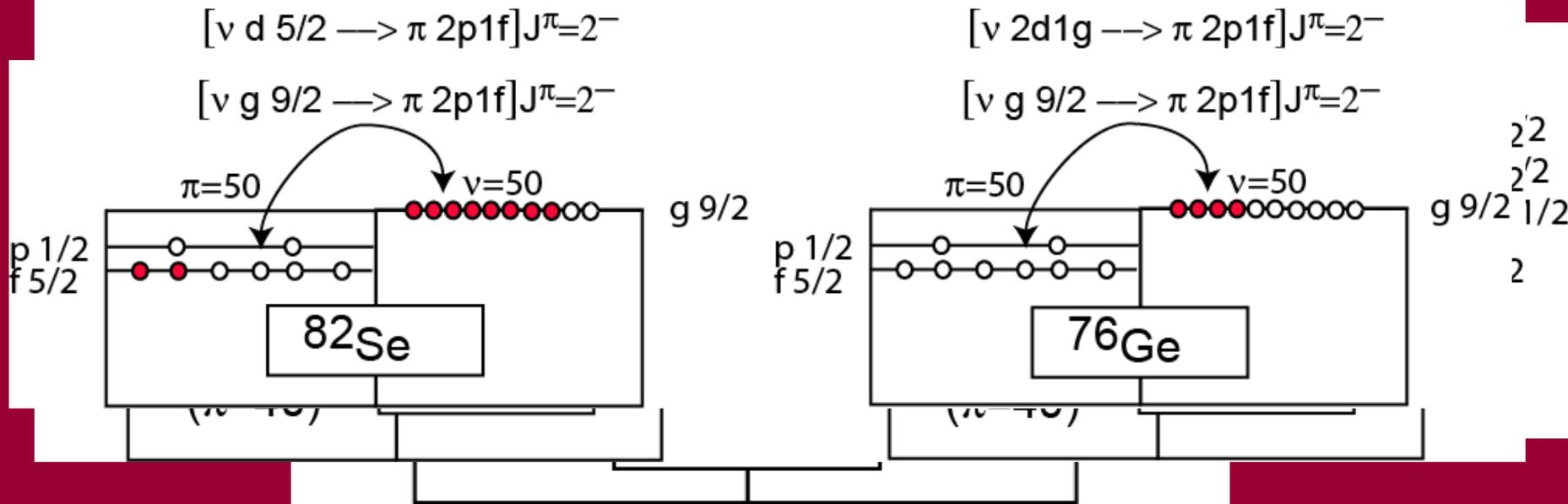
The 2^- strength makes up
 ~ 20-30% of the $0\nu\beta\beta$ ME!!

J. Suhonen, Phys. Lett B607, 87 (2005)

Expmt:

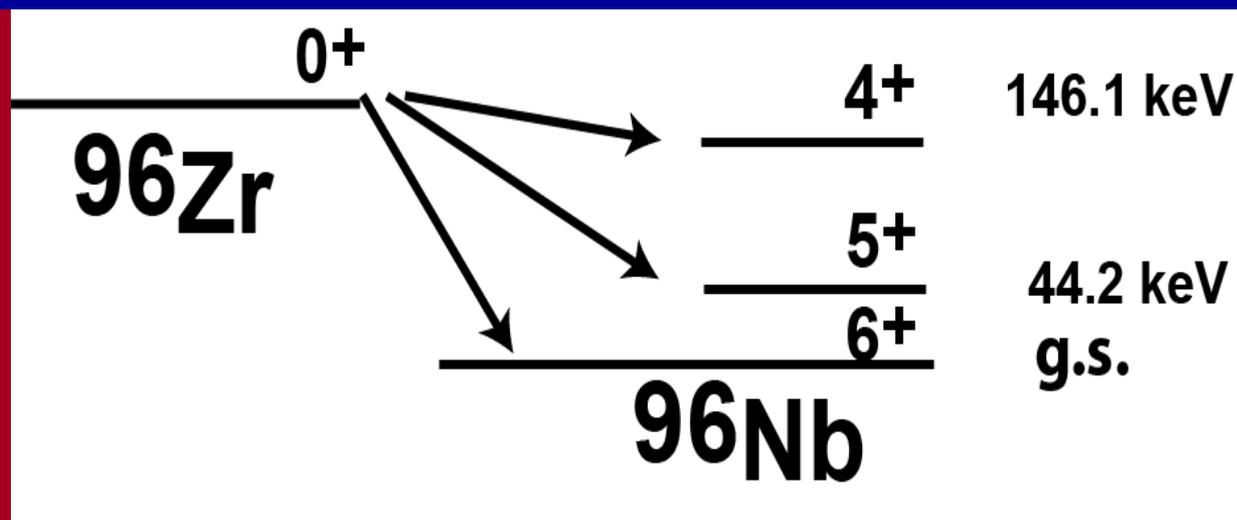
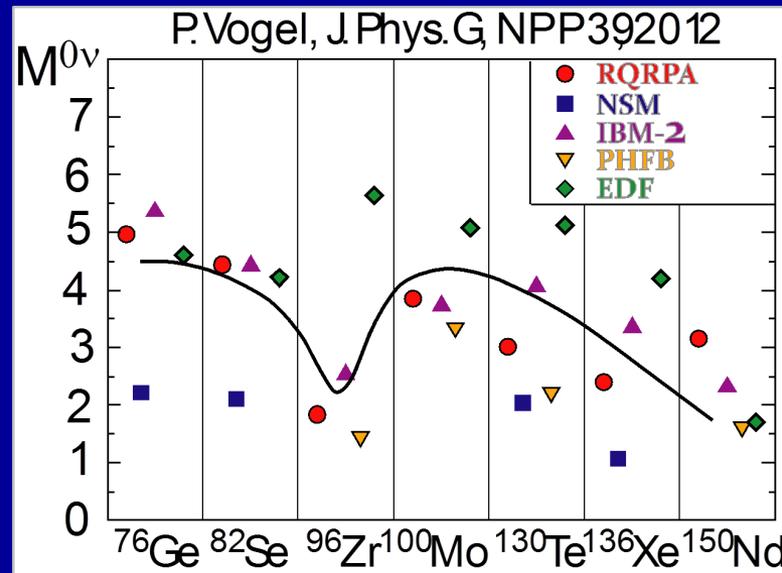
^{136}Xe exhibits largest 2^- strength
 $0\nu\beta\beta$ ME enhanced???





3rd piece of puzzle

The A=96 system and the $0\nu/2\nu$ $\beta\beta$ NMEs



Single β & $\beta\beta$ decay in ^{96}Zr

two conflicting half-lives:

NEMO-3: $T_{1/2}^{2\nu\beta\beta} = (2.3 \pm 0.2) \times 10^{19} \text{ y}$
 geo-chem: $T_{1/2}^{\beta} = (0.94 \pm 0.32) \times 10^{19} \text{ y}$ ①

Q: can this difference be reconciled ?

A: yes, if single β competes with $\beta\beta$ decay

$$(T_{1/2})^{-1} = (T_{1/2}^{2\nu\beta\beta})^{-1} + (T_{1/2}^{\beta})^{-1}$$

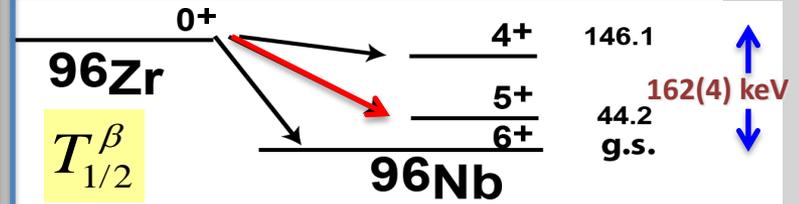
expected $T_{1/2}^{\beta} = (1.6 \pm 0.9) \times 10^{19} \text{ y}$

experiment $T_{1/2}^{\beta} > 2.6 \times 10^{19} \text{ y}$ ②

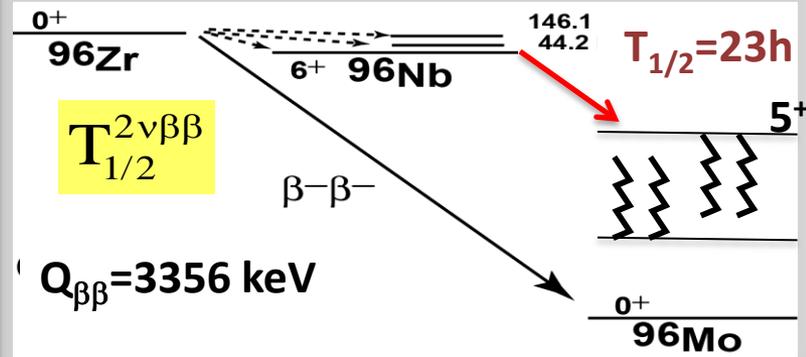
pred. (QRPA) $T_{1/2}^{\beta} = 24 \times 10^{19} \text{ y}$ ③

BUT

$$(T_{1/2}^{\beta})^{-1} \propto o(Q^{13}) \langle M_{\beta}^{4u} \rangle^2$$



$0^+ \rightarrow 6^+$ 6-fold non-unique (unobservably long)
 $0^+ \rightarrow 5^+$ 4-fold unique (possible)
 $0^+ \rightarrow 4^+$ 4-fold non-unique (no phase space)

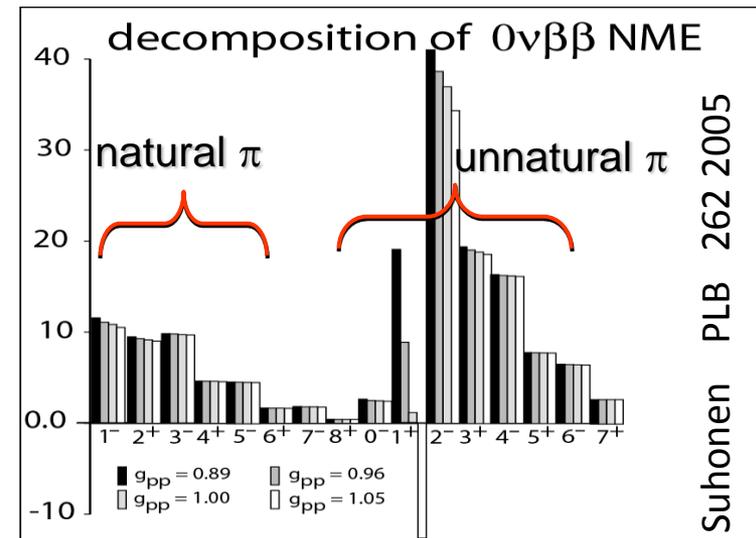
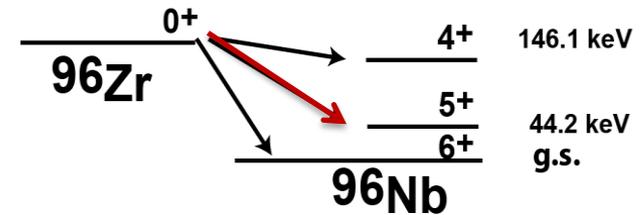


Q-value

$$\longrightarrow M_{\beta}^{4u} \longrightarrow (T_{1/2}^{0\nu\beta\beta})^{-1} \propto Q^5 |M_{\beta\beta}^{0\nu}|^2 \langle m_{\beta\beta} \rangle^2$$

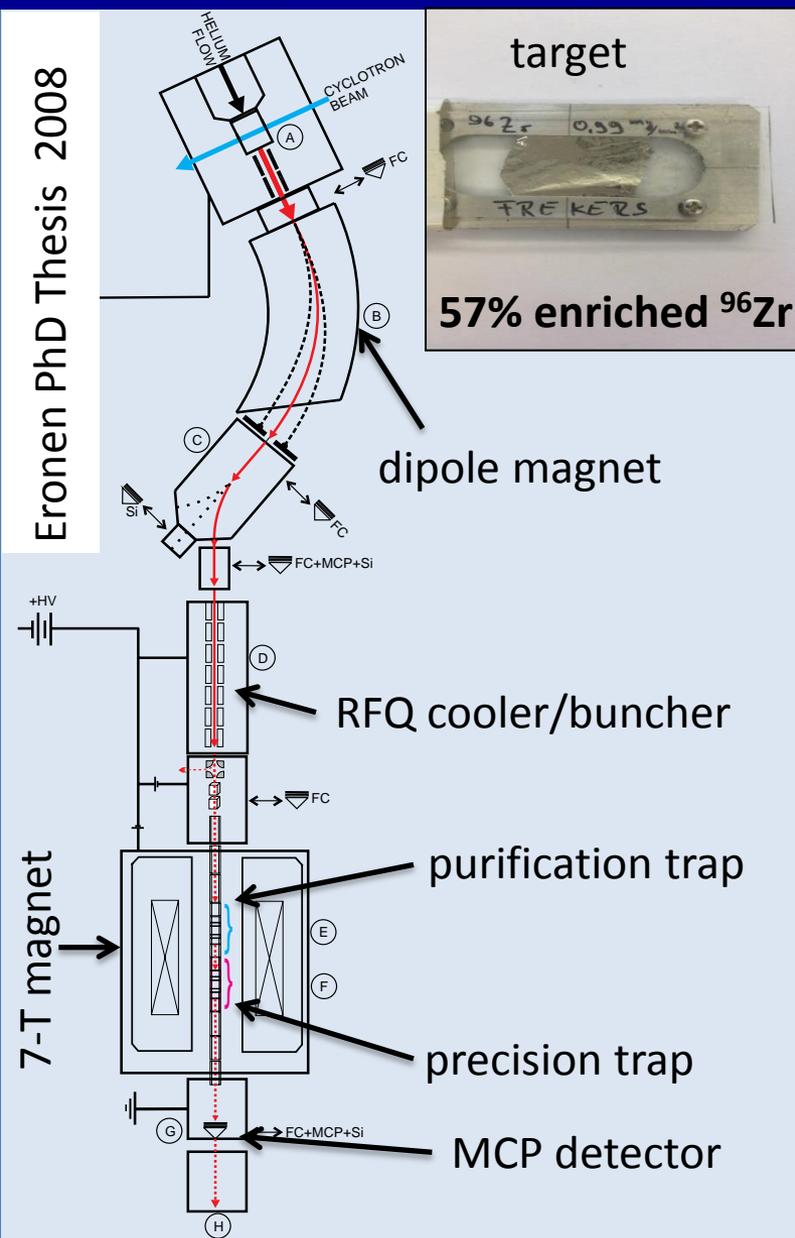
Idea

- (i) measure **Q-value** for $^{96}\text{Zr} \longrightarrow ^{96}\text{Nb}$ **single β -decay** by precision mass measurement and
- (ii) measure the **single β -decay** rate
- (iii) \rightarrow ft-value
- determine the ^{96}Zr **4-fold forbidden β -decay NME** and confront with theory
- confront with same theories aimed at calculating **$0\nu\beta\beta$ -decay NME for the same nucleus!!**



IGISOL/JYFLTRAP mass measm'nts

Eronen PhD Thesis 2008



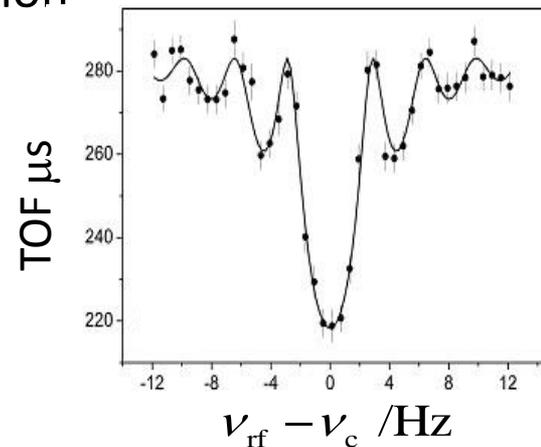
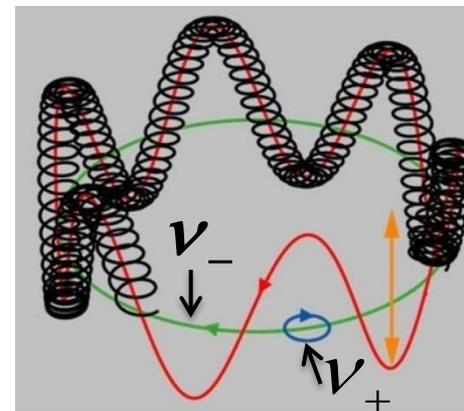
**^{96}Zr (p, n) ^{96}Nb reaction
for production of ^{96}Nb**

- performing accurate mass measurements via **cyclotron frequency**

$$v_c = \frac{1}{2\pi} \frac{q}{m} \cdot B$$

$$v_c = v_- + v_+$$

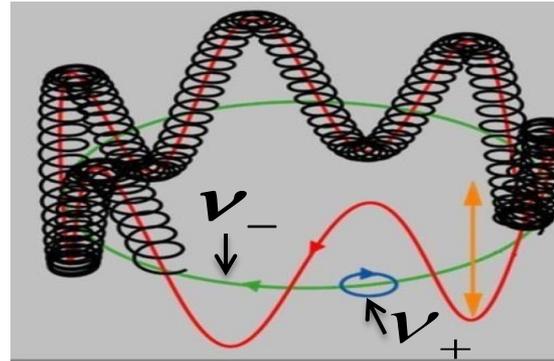
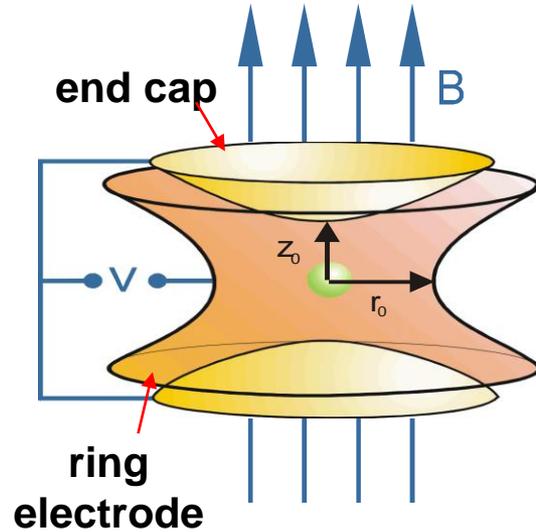
- frequency determination done by **TOF-ICR** technique



ion motion in a Penning Trap

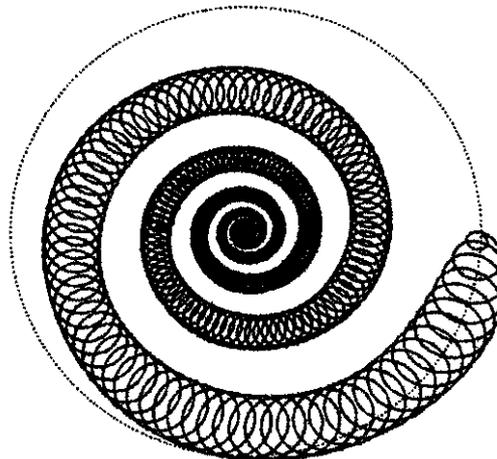
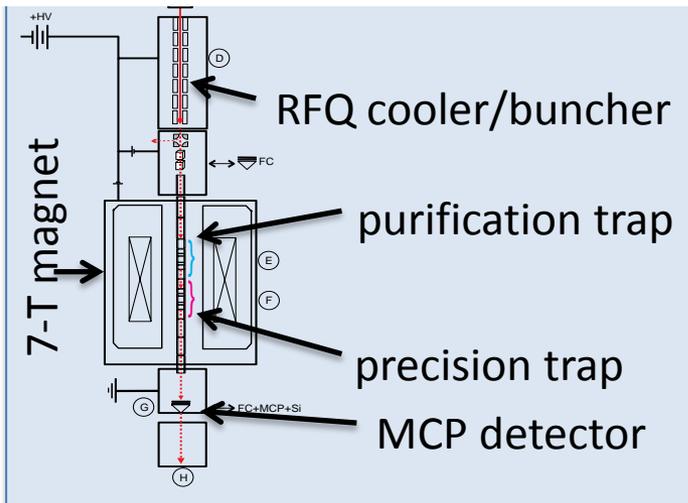
a superposition of 3 harmonic motions:

- 1) axial motion (ν_z)
- 2) magnetron motion (ν_-)
- 3) reduced cyclotron motion (ν_+)



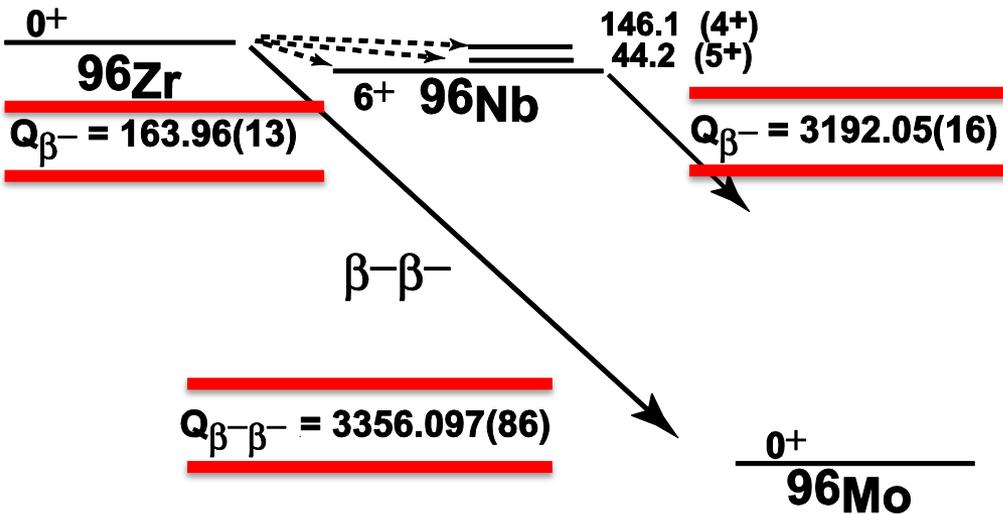
typical freq's: $\nu_- \approx 1$ kHz,
 $\nu_+ \approx 1$ MHz
 cycl. freq.: $\nu_C = \nu_- + \nu_+$

principle of isobar separation in purification trap



excitation of ion of interest with cyc. freq. causes collisional cooling of this species
 (purification: $> 500/1$)

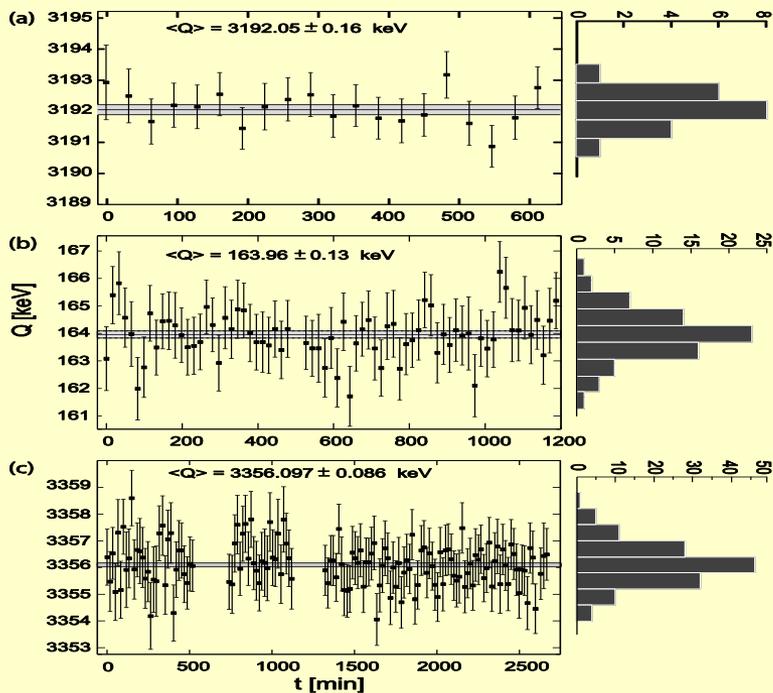
Results



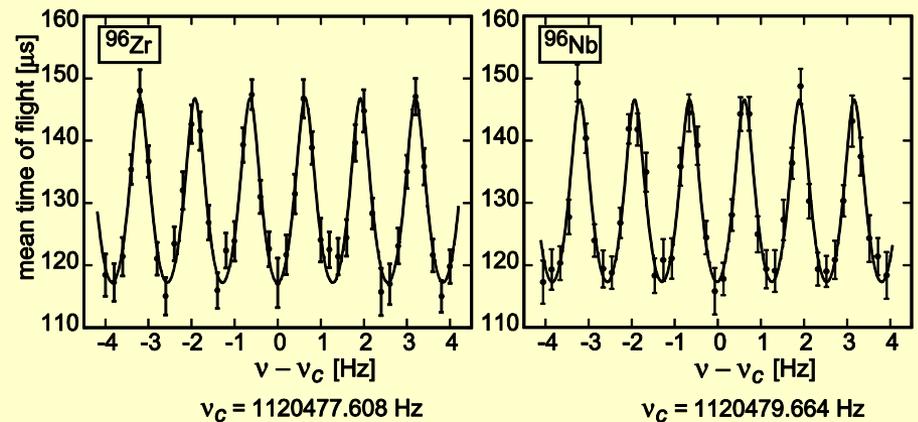
^{96}Zr

$Q_{\beta\beta} = 3356.097 \pm 0.086 \text{ keV}$
 7.1 keV higher than AME2012

$Q_{\beta} = 163.96 \pm 0.13 \text{ keV} \quad !!$



Ramsey excitation



Next: need $T_{1/2}$ of single β decay

$$T_{1/2}(\text{QRPA}) = 24 \times 10^{19} \text{ yr} \quad (g_A = 1)$$

$$T_{1/2}(\text{SM}) = \frac{11}{g_A^2} \times 10^{19} \text{ yr} \quad (g_A^2 = 0.75 \dots 1.6)$$

$$T_{1/2}(\text{exp}) > 2.3 \times 10^{19} \text{ yr}$$

Important side effect:

single β decay depends on g_A^2
2 ν /0 $\nu\beta\beta$ decay depends on g_A^4

A measurements of single β decay would give for the first time an expmtl handle on the quenching of the axial vector coupling constant³⁶

$T_{1/2}$ of ^{96}Zr decay

Geochemical method (zircon sample)

accumulation of decay daughter over geological time

- **Pros:** can provide longer $T_{1/2}$
- **Cons:** - measure only total decay
- systematics from ^{95}Mo (n, γ) reaction
n-sources: ^{238}U sf, cosmics
BUT: zircon contains Gd (i.e. $^{157}\text{Gd}(n, \gamma)$)
- needs chemical isobar separation ($A=96$)

age of sample: $1.822(\pm 0.003) \times 10^9$ y
dated by U-Pb technique
origin: Capel Australia
advantage: Mo is embedded in zircon
attice with infinite retention time

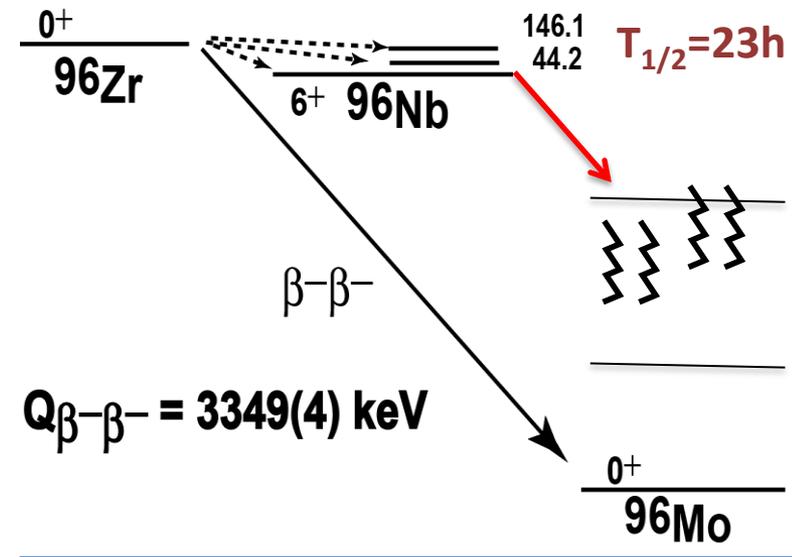
geo. $T_{1/2}$ calculation

- # of ^{96}Mo excess atoms in sample
- # of ^{96}Zr atoms in sample
- age of the sample

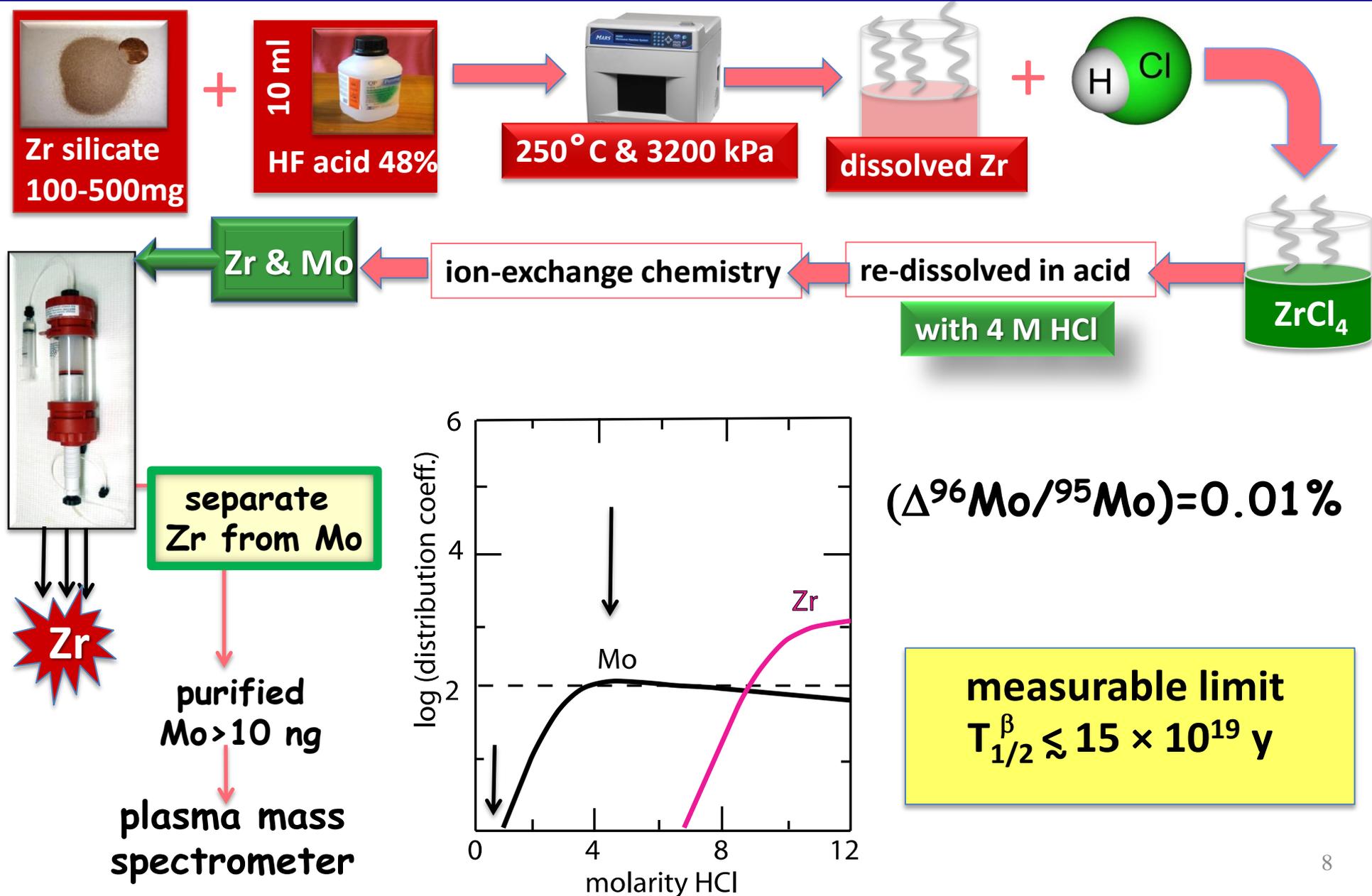
Counting-rates method

measures the decay rate

- **Pros:** measure partial decays
- **Cons:** - extremely low counting rates
- background and time limited



Isobar separation method (U of Calgary)



These and more pieces to the puzzle

- **chargex-reactions for $2\nu\beta\beta$ decay**
Hadronic chargex and weak-interaction x-sections are fortuitously connected -- exploit this!!
- **$0\nu\beta\beta$ -decay and chargex-reaction**
useful but limited
e.g. 2^- states (resol'n is key)
- **mass measurements for $\beta\beta$ -decay NME**
the potential still needs to be exploited
- **single β decay half-life for ^{96}Zr (and ^{48}Ca)**
- **ground-state properties of intermediate odd-odd nuclei (TRIUMF, TITAN-EC)**
- **need to address quenching issue urgently!!**
- **need relevant spectroscopic information**
(but theory needs to tell what is relevant)
- **theory needs to converge !!**
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